

Basics about cellular biology

Stéphane Avril

IMWS

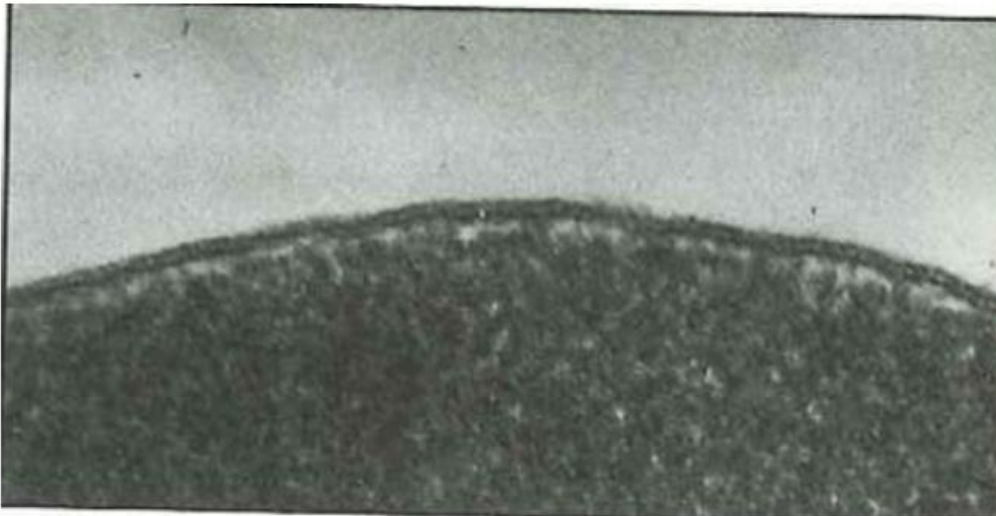
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The cell membrane

Envelope that acts as a selective barrier between cytoplasm and extracellular medium.

Without plasmic membrane, the cell could not maintain its ordered chemical system integrity.



Cell membrane

5 or 6 nm thick

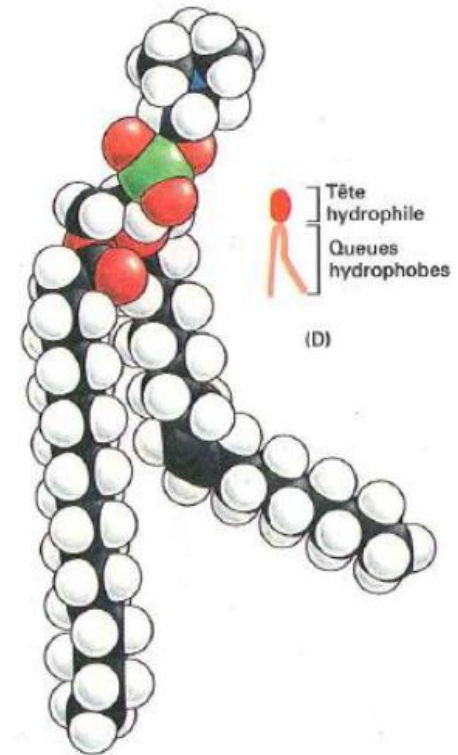
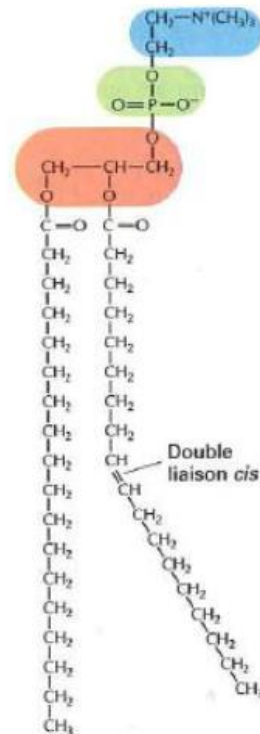
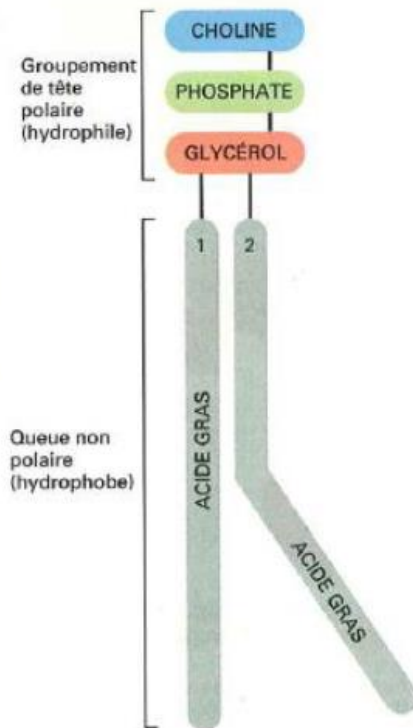
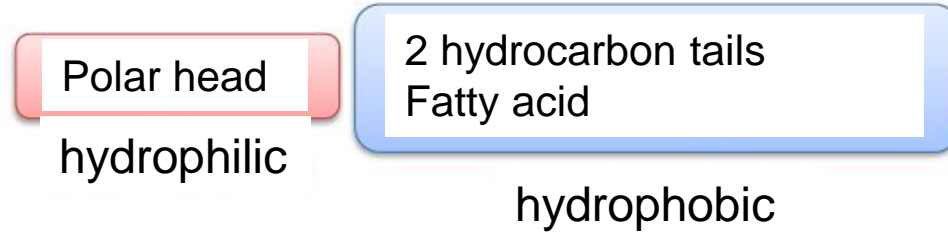
Composition: lipids and proteins

The lipid bilayer

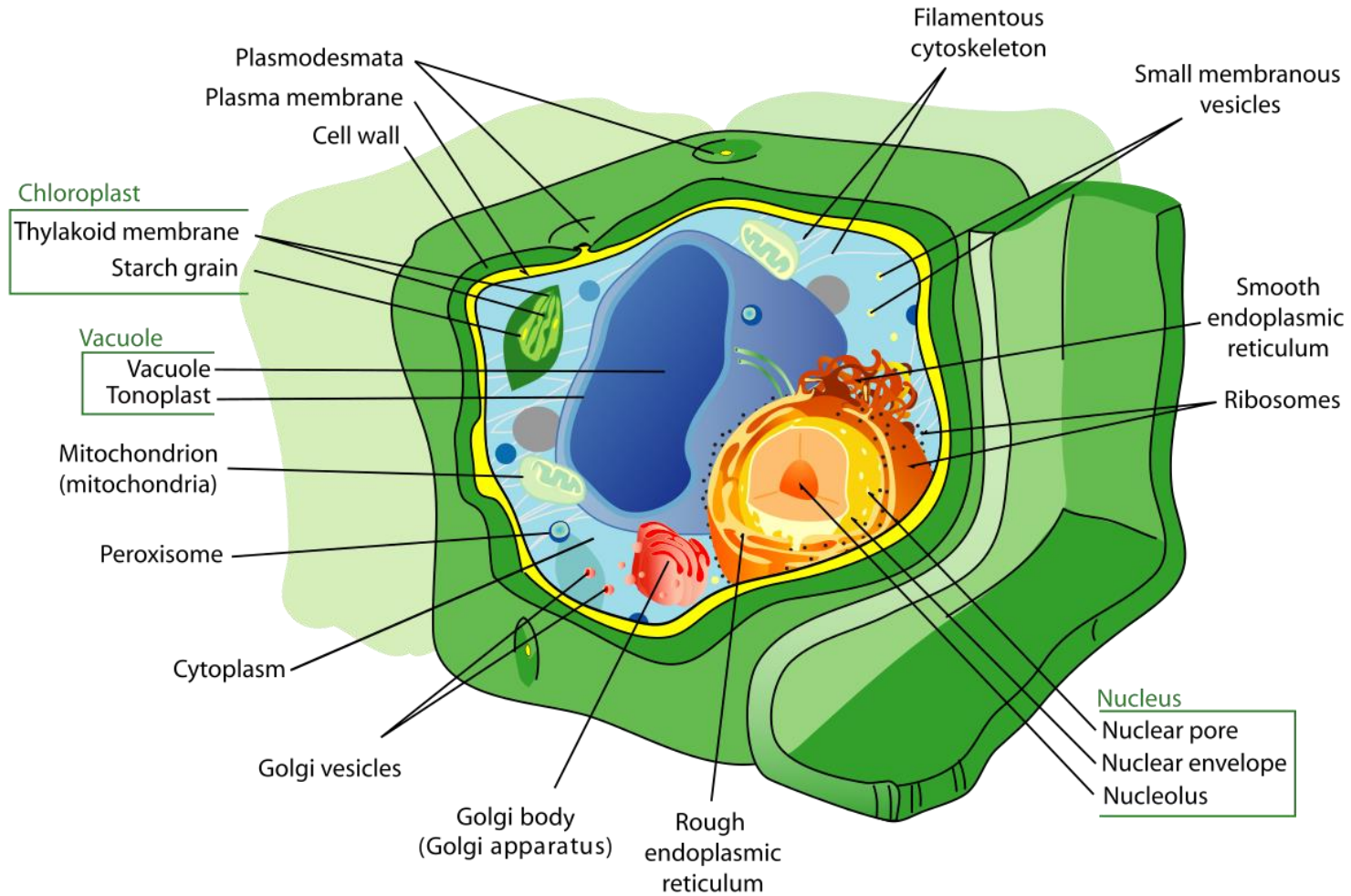
3 main types of membrane lipids:

- Phospholipids (the most numerous)
- Cholesterol
- Glycolipids

Structure of lipids :



Structure of plant eukaryote cell



Plastids are found in plants and algae.

The best known are chloroplasts, in the cells of photosynthetic organisms, which convert light energy into chemical energy used to make sugars from carbon dioxide.

They also have their own genome.

In plants, algae and fungi, the cell is surrounded by a pectocellulosic cell wall which provides the body with a skeleton. Deposits of compounds such as **suberin** or **lignin** modulate the physicochemical properties of the wall, making it more solid or more impermeable.

Layout

I. Definition and general presentation of the cell.

II. The main cellular structures.

III. The origin of cells.

IV. The different cellular organizations.

V. Cellular homeostasis.

VI. Structure of the eukaryotic cell:

1. Animal:

a. The organelles.

b. The membrane.

2. Plants.

VII. Genetic information.

Genetic information

DNA = Deoxyribonucleic acid

phosphate

base

Sugar
desoxyribose

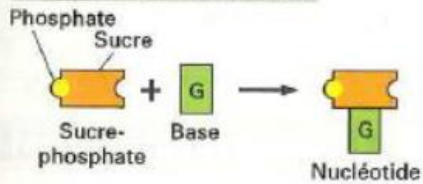
A= Adénine
T= Thymine
C= Cytosine
G= Guanine

Unité de
base

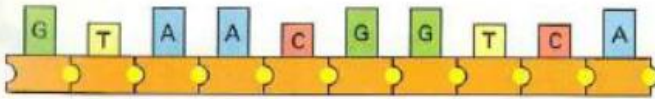
Brin
d'ADN

Double
brin
d'ADN

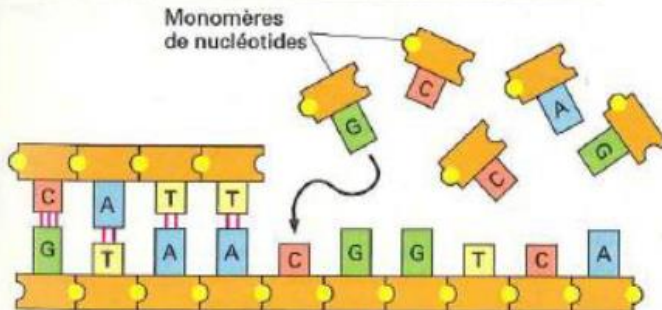
(A) Unités de structure de l'ADN



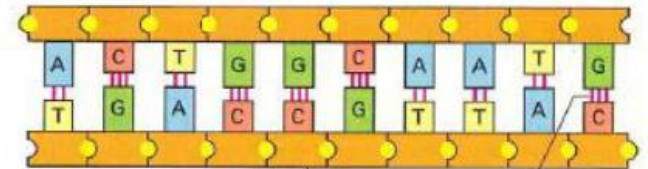
(B) Brin d'ADN



(C) Polymérisation à l'aide d'une matrice d'un nouveau brin



(D) ADN double brin



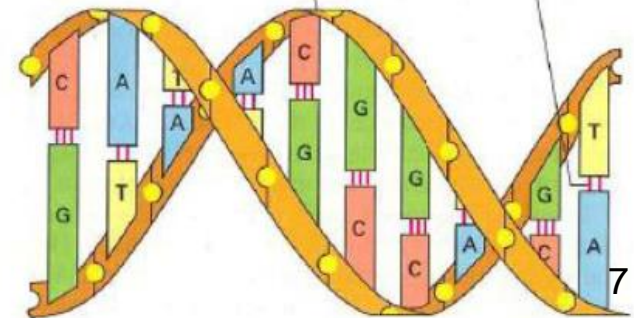
Squelette
sucre-phosphate

Liaisons hydrogène entre
les paires de bases

Double
brin
d'ADN

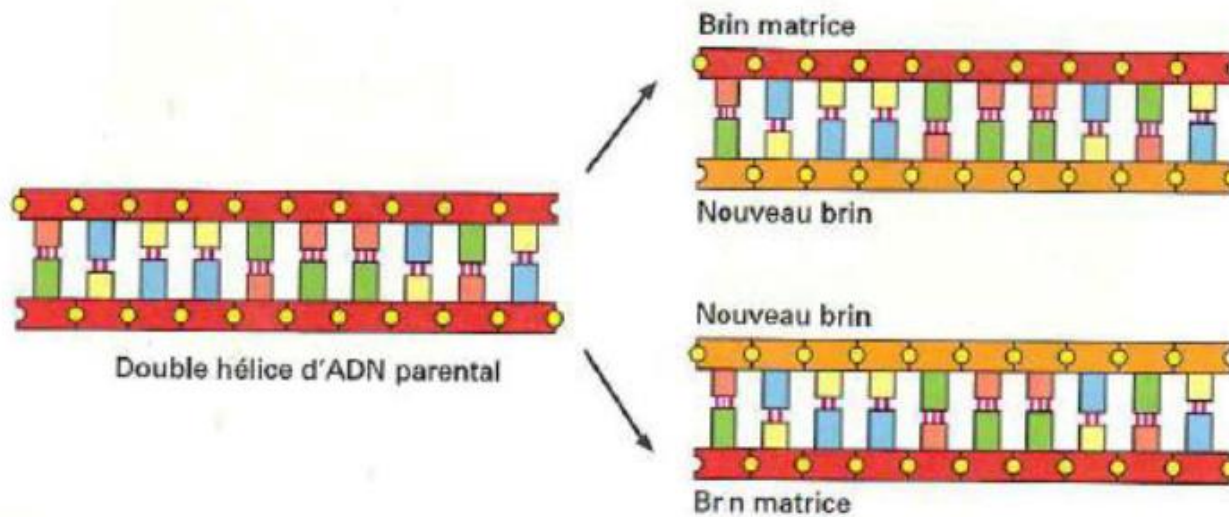
Double
hélice
d'ADN

(E) Double hélice d'ADN

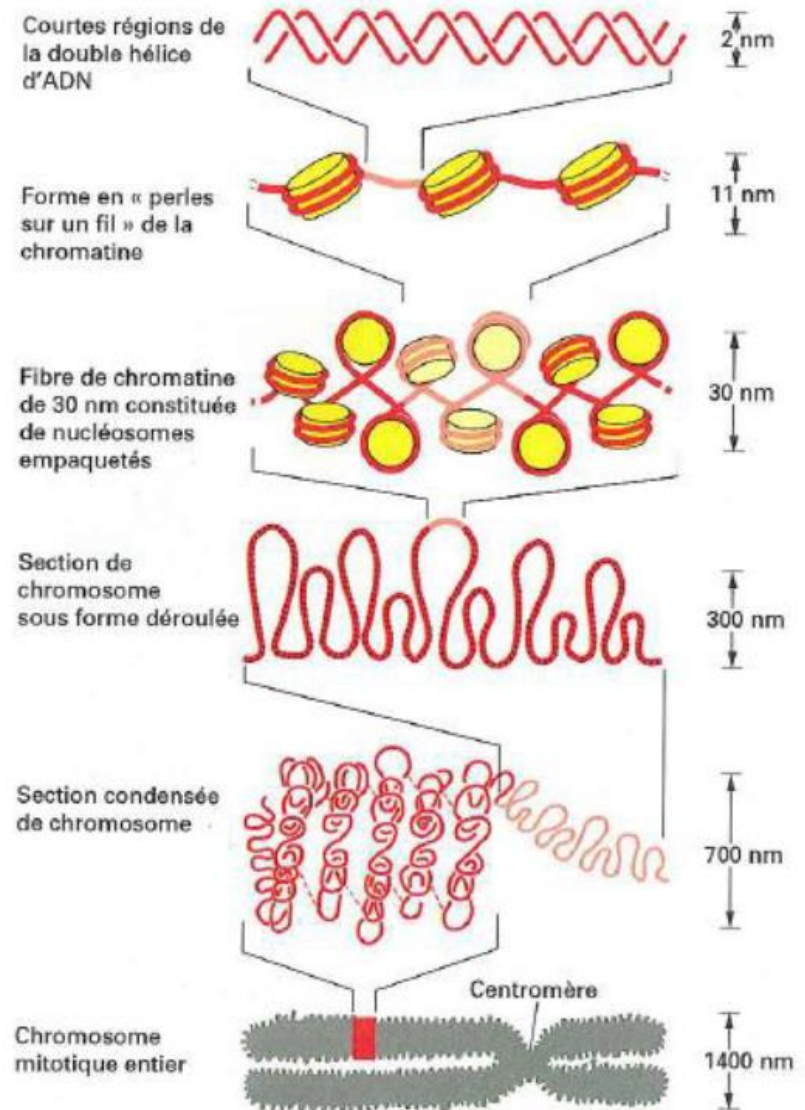


- Duplication of genetic information by repping DNA:

→ polymerization using a matrix.



Each cell contains 2 m of DNA. The nucleus measures 6 μm diameter
→ compaction of the DNA with proteins



23 pairs of chromosomes

22 pairs of homologous chromosomes

1 pair of sexual chromosome XX or XY

RÉSULTAT NET : CHAQUE MOLÉCULE D'ADN A ÉTÉ EMPAQUETÉE DANS UN CHROMOSOME MITOTIQUE QUI EST 10 000 FOIS PLUS COURT QUE SA LONGUEUR DÉROULÉE

Cellular mechanics

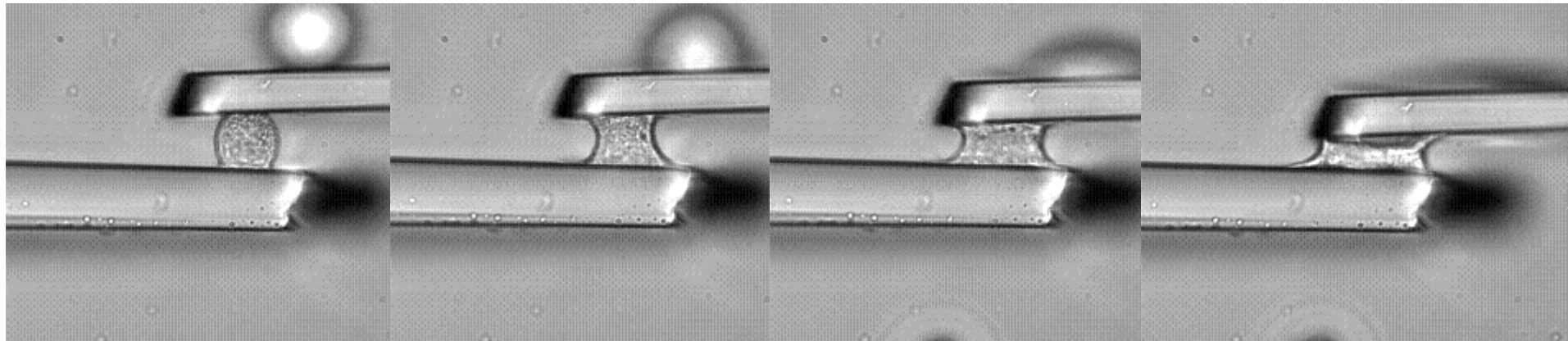
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Passive and active mechanical properties of isolated living cells



Cell and mechanics

Passive properties

Viscoelasticity of tissues

structure \leftrightarrow propagation of forces

Active properties

Mechanical response to a biochemical signal

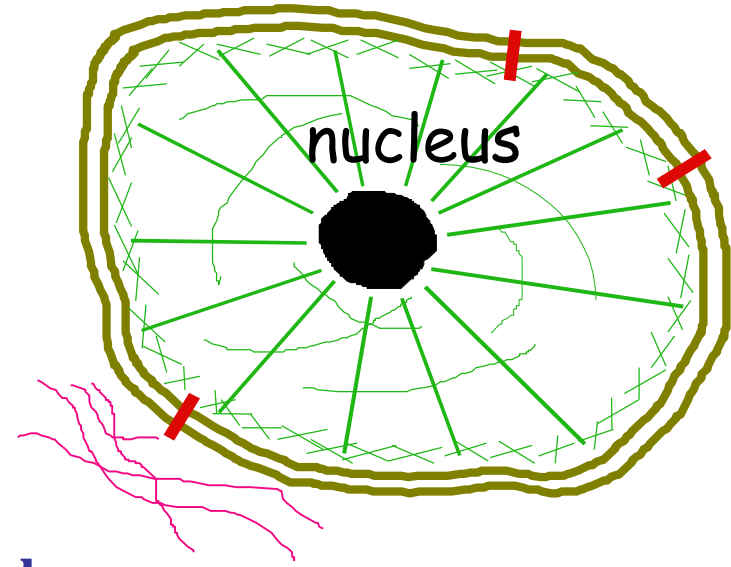
chemo-attractant, adhesion \rightarrow polarisation, migration

Mechanical / biochemical response to a mechanical signal :

mechanotransduction

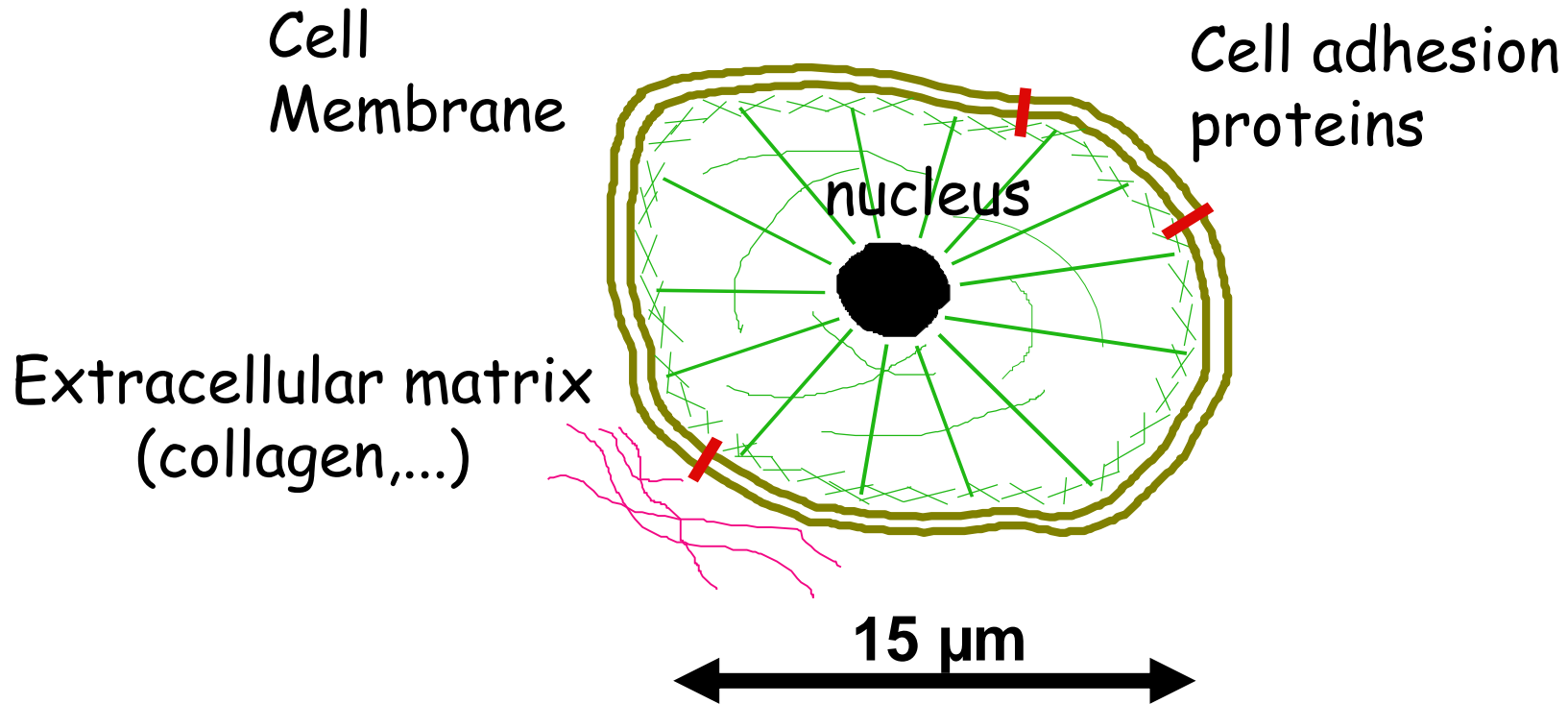
cultured cells: always under tension

secretions of an epithelium modulated by blood flow



Mechanics \leftrightarrow **Biological function**

Cell and mechanics



Dynamic cytoskeleton (2D-3D) :

- Actin filaments ($\text{Ø}=8\text{nm}$, $L_p=15\mu\text{m}$)
- Intermediate filaments ($\text{Ø}\approx 10\text{nm}$, $L_p\approx 500\text{nm}$)
- Microtubules ($\text{Ø}=25\text{nm}$, $L_p \approx \text{qq mm}$)

Passive properties

Apply a **controlled stress** to the cell and determine its :

Mechanical response

Strain measurement (creep)

Assessment of the viscoelastic modulus

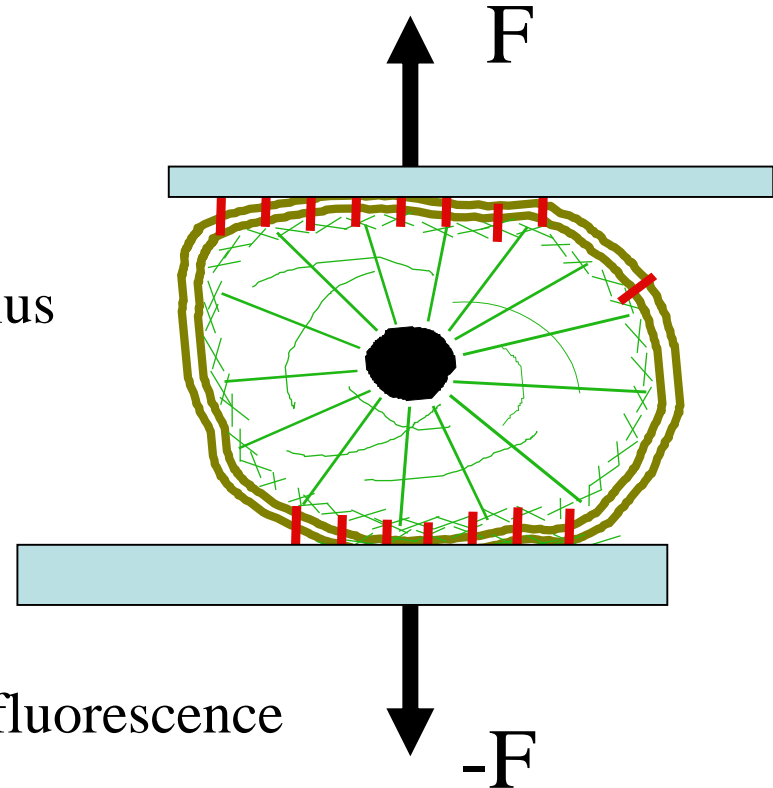
With collaborations

Structural response

visualization of the cytoskeleton by fluorescence

Biochemical response (genetics)

monitoring of protein or genetic markers



Active properties

Observe the evolution of a cell in a predefined and simple geometry (3D):

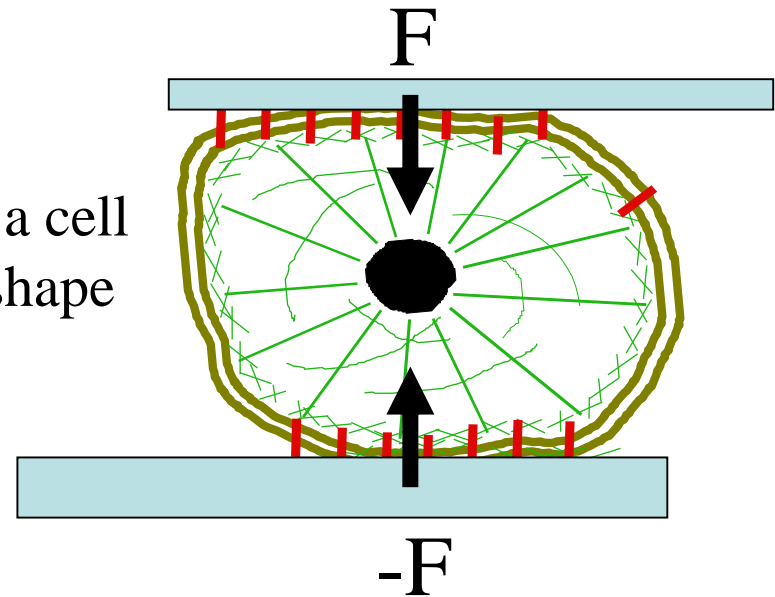
Mechanical activity

Measurement of the force applied by a cell
Correlation with the evolution of its shape

With collaborations

Structural response

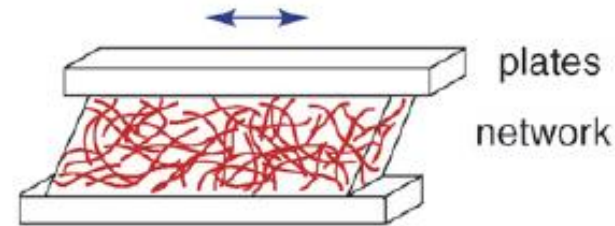
visualization of the cytoskeleton by fluorescence



Characterization methods

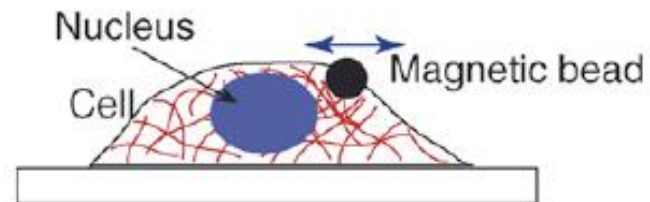
Bulk rheology

A material is sheared between two plates using an oscillatory stress to probe the shear elastic, G' , (in-phase) and viscous, G'' , (out-of-phase) moduli.



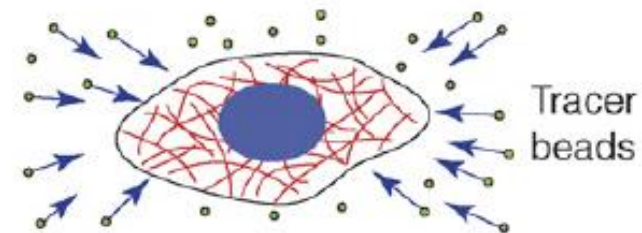
Magnetic bead cytometry

An external magnetic field applies a stress to a magnetic bead. The bead's position is tracked to determine the response.



Traction force microscopy

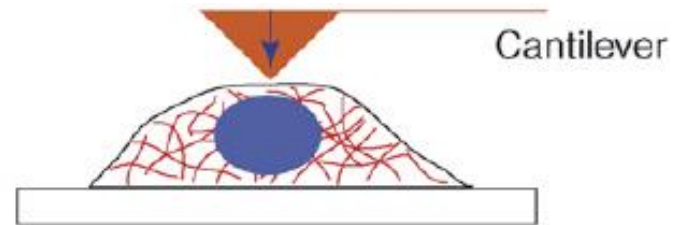
Cell contractions deform a flexible substrate. Forces are estimated from bead displacements.



Characterization methods

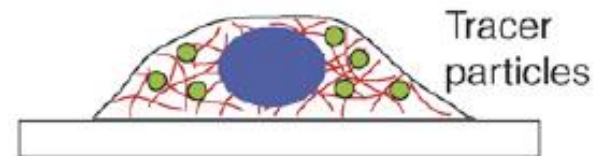
Atomic force microscopy

A cantilever applies stress to the cell. The cantilever deflection is measured by laser reflection.



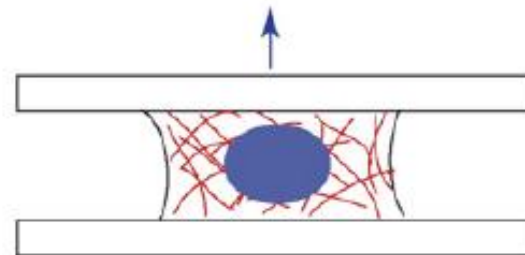
Microrheology

The motion of probe particles is measured using either video or laser tracking techniques. Particle motion is either driven externally or thermally induced and is interpreted to yield the viscoelastic modulus.



Whole cell stretching

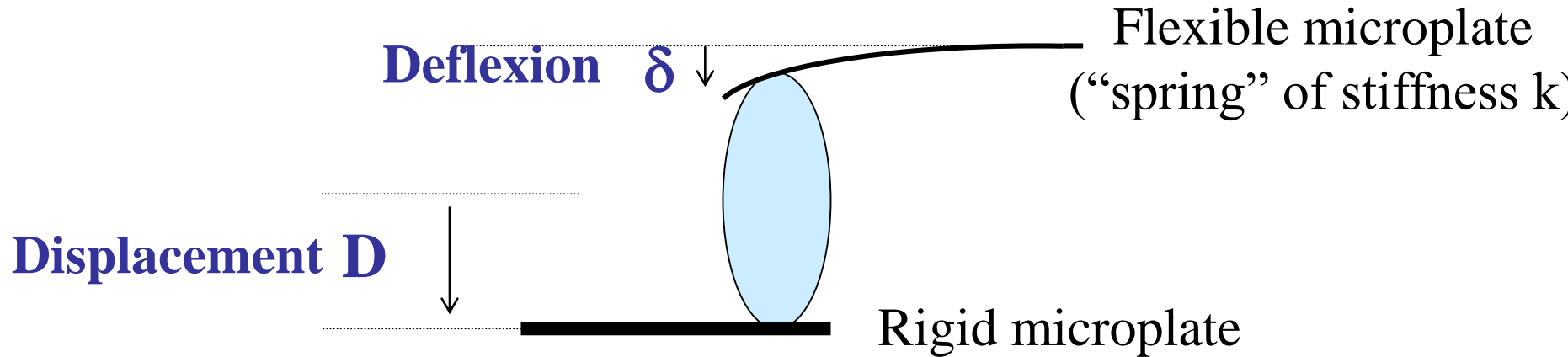
A cell is attached to two surfaces. A force is applied to one surface and the plate separation is measured.



Uniaxial stretching



Uniaxial stretching



Force

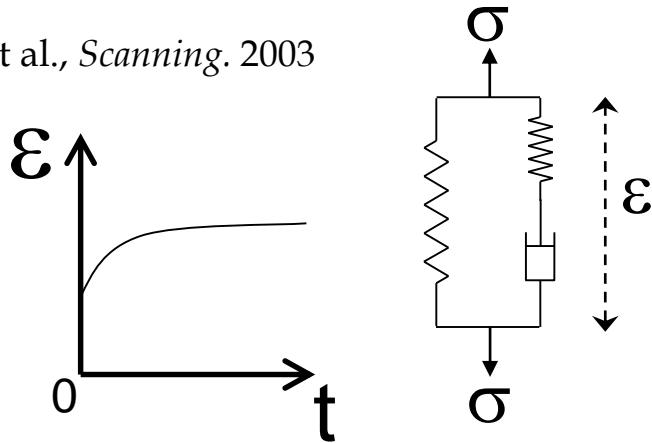
$$F = k \delta$$

Passive properties:

Rheology of a living cell using uniaxial stretching

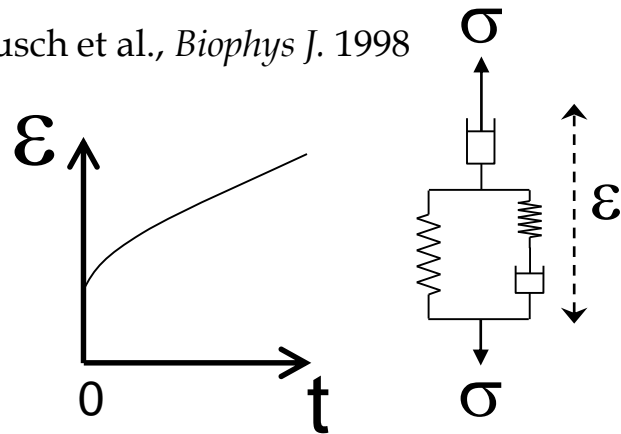
Local measurements, time domain

Wu et al., *Scanning*. 2003



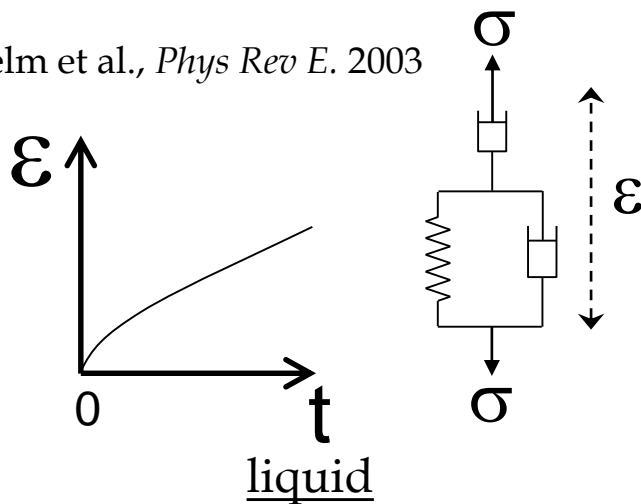
Standard linear solid

Bausch et al., *Biophys J*. 1998



Liquid with instantaneous elasticity

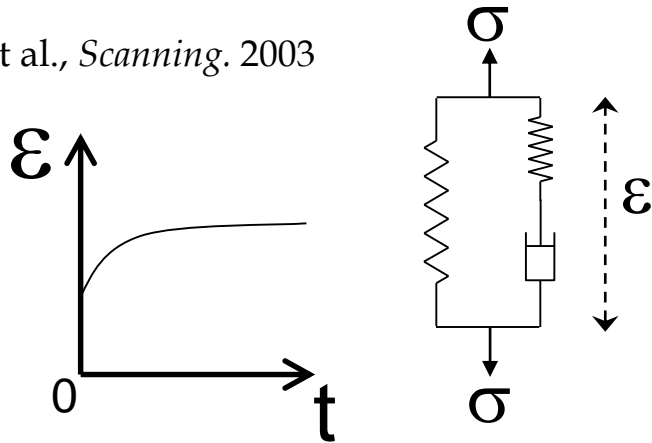
Wilhelm et al., *Phys Rev E*. 2003



liquid

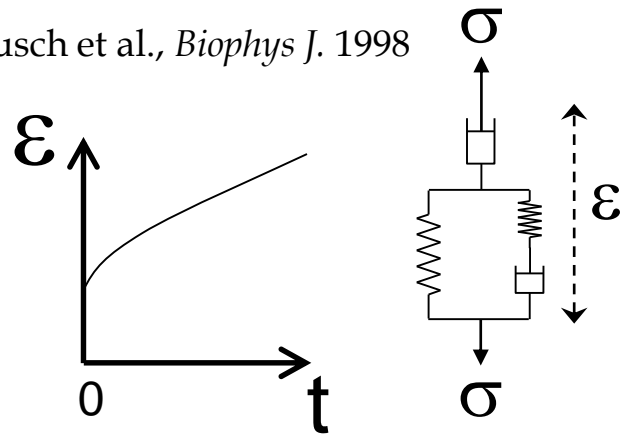
Local measurements, time domain

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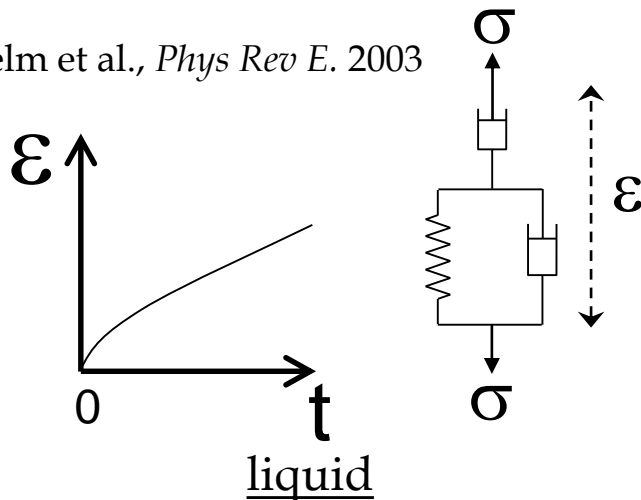
Standard linear solid

Bausch et al., *Biophys J*. 1998



Liquid with instantaneous elasticity

Wilhelm et al., *Phys Rev E*. 2003



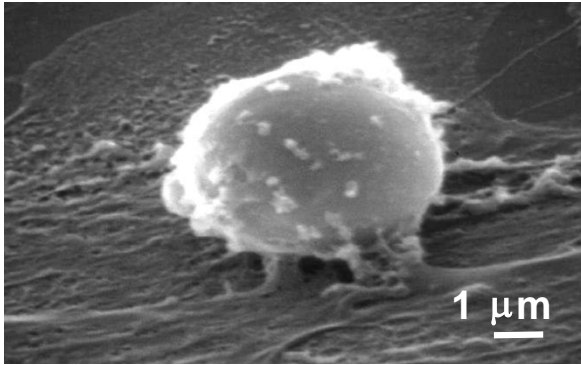
liquid

Some characteristic times

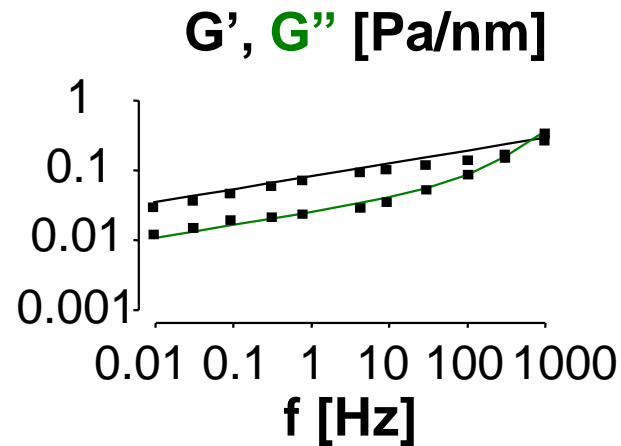
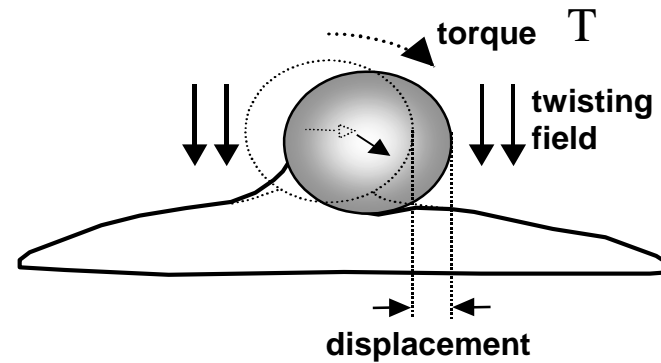
Viscous dissipation

Very different behaviors

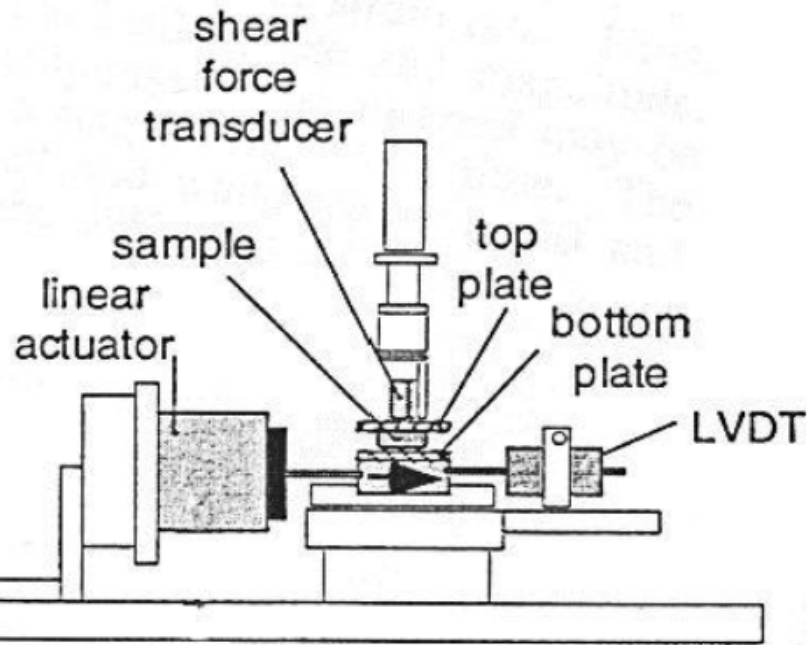
Local rheometry, frequency analysis



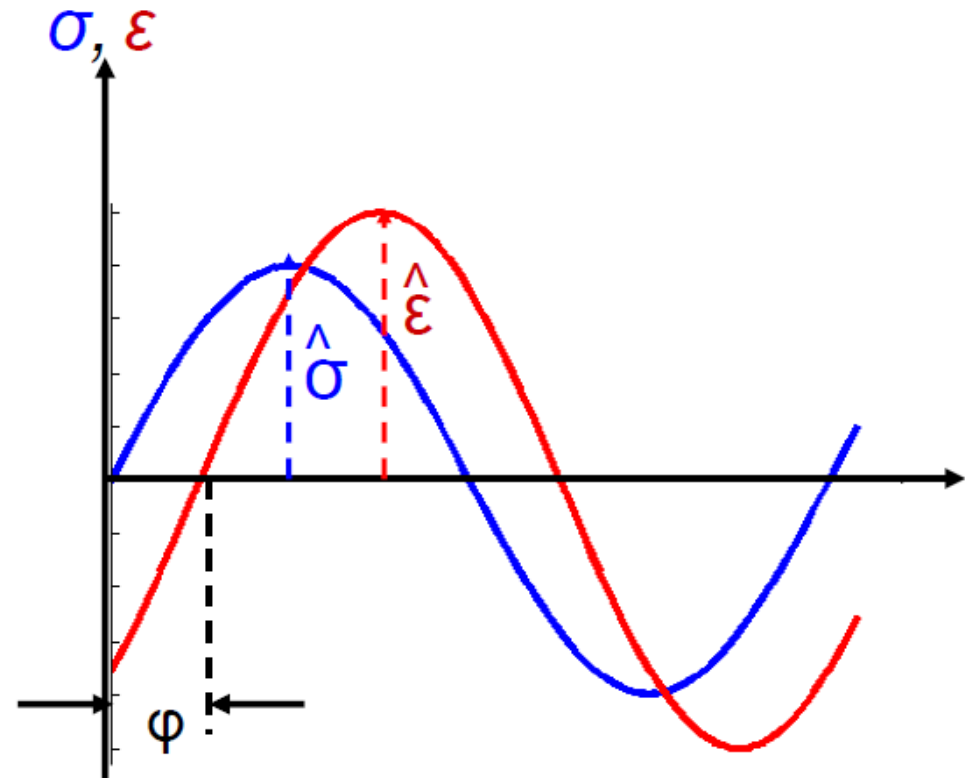
Fabry et al., *Phys Rev Lett*. 2001



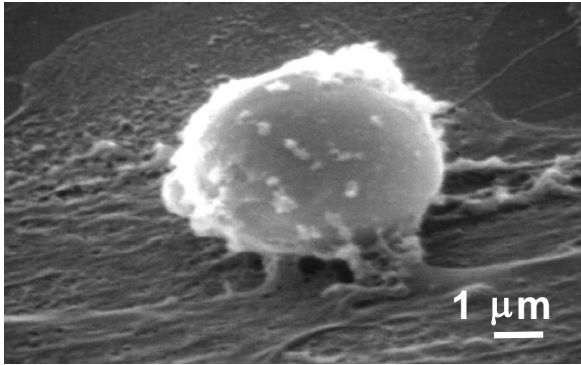
Oscillations to determine viscoelastic properties



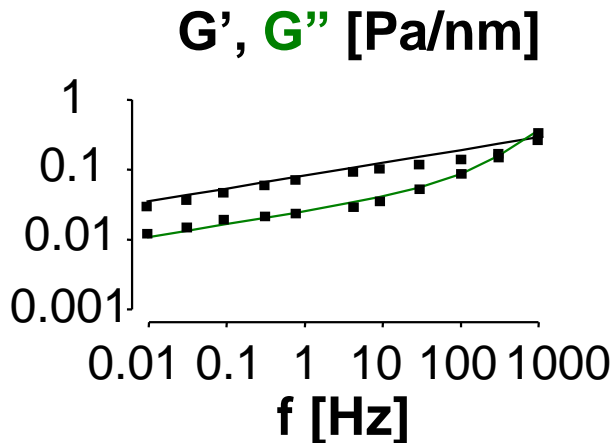
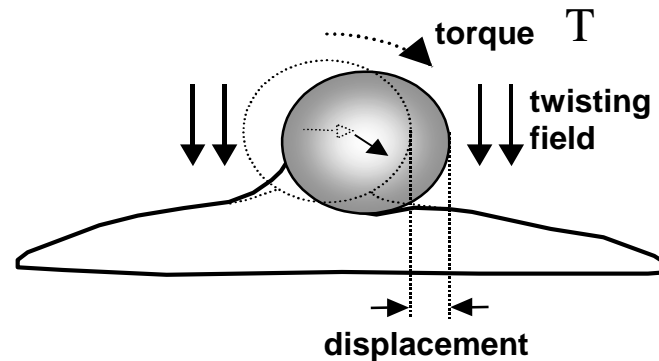
$$\varepsilon = \hat{\varepsilon} \sin(\omega t)$$



Local rheometry, frequency analysis



Fabry et al., *Phys Rev Lett*. 2001



$$G^*(\omega) = G_0 \Gamma(2-x) (j \omega \tau_0)^{x-1} + j \omega \mu$$

$$\Gamma : z \mapsto \int_0^{+\infty} t^{z-1} e^{-t} dt \quad \Gamma(z+1) = z \Gamma(z).$$

No characteristic time

Elasticity and dissipation from same origin

Unique behavior preserved

AVERAGE !

Soft glassy medium behavior

Out of balance

Structural disorder

Metastability

Effective temperature (glass transition)

What if the curve of the model does not fit the curve of the material we want to describe?

Generalized Models...

- The strain response to an arbitrary stress history is obtained from $J(t)$ by *superposition*

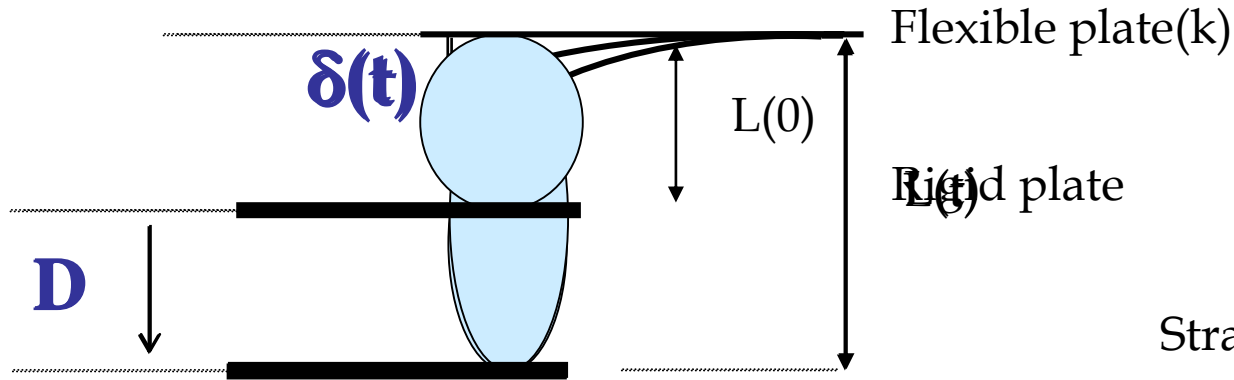
$$\varepsilon(t) = \int_0^t J(t - \tau) d\sigma = \int_0^t J(t - \tau) \frac{d\sigma}{d\tau} d\tau$$

- The strain response to an arbitrary stress history is obtained from $J(t)$ by *superposition*

$$\varepsilon(t) = \int_0^t J(t - \tau) d\sigma = \int_0^t J(t - \tau) \frac{d\sigma}{d\tau} d\tau$$

1. From uniaxial stretching to single cell rheometer

Uniaxial stretching ($\dot{\sigma} \neq 0$ et $\dot{\varepsilon} \neq 0$)



Stress :

$$\sigma(t) = \frac{F(t)}{S}$$

$$F(t) = k \cdot \delta(t)$$

Strain :

$$\varepsilon(t) = \frac{L(t) - L(0)}{L(0)}$$

Stress-strain relationship

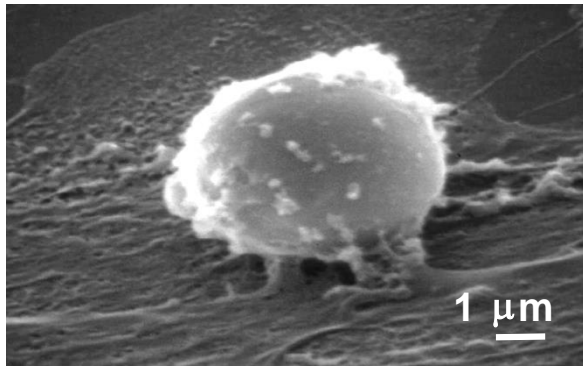
$$\varepsilon(t) = J(t)\sigma(0) + \int_0^{+\infty} J(t-t')\dot{\sigma}(t')dt'$$

$\dot{\sigma} \neq 0 \Rightarrow$ Very difficult to determine J

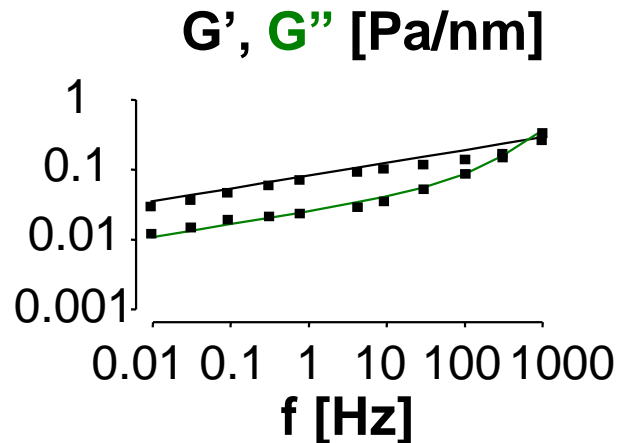
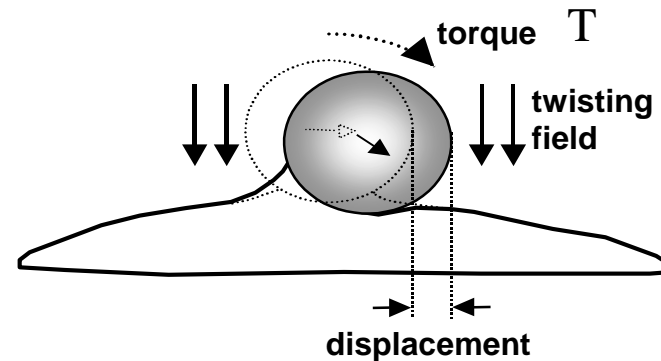
Avoid convolution product \Leftrightarrow oscillations ($\sigma(\omega)$)

ou constant stress ($\dot{\sigma} = 0$)

Local rheometry, frequency analysis



Fabry et al., *Phys Rev Lett*. 2001



$$G^*(\omega) = G_0 \Gamma(2-x) (j \omega \tau_0)^{x-1} + j \omega \mu$$

No characteristic time
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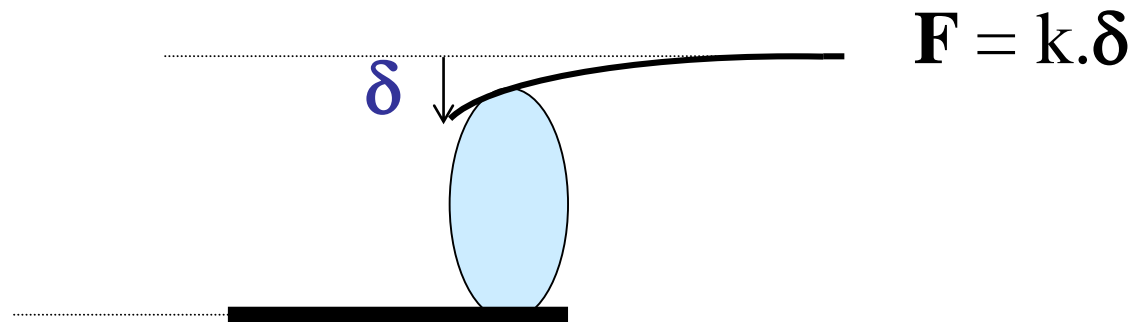
Effective temperature (glass transition)

1. From uniaxial stretching to single cell rheometer

Rheometer ($\dot{\sigma} = 0$)

$$\varepsilon(t) = J(t)\sigma(0)$$

At constant stress: measurement of $J \Leftrightarrow$ measurement of strain ε



$$\dot{\sigma} = 0 \Leftrightarrow \delta \text{ constant}$$

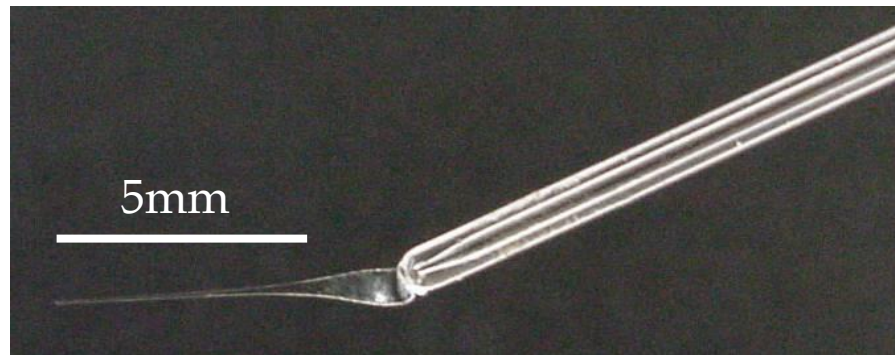
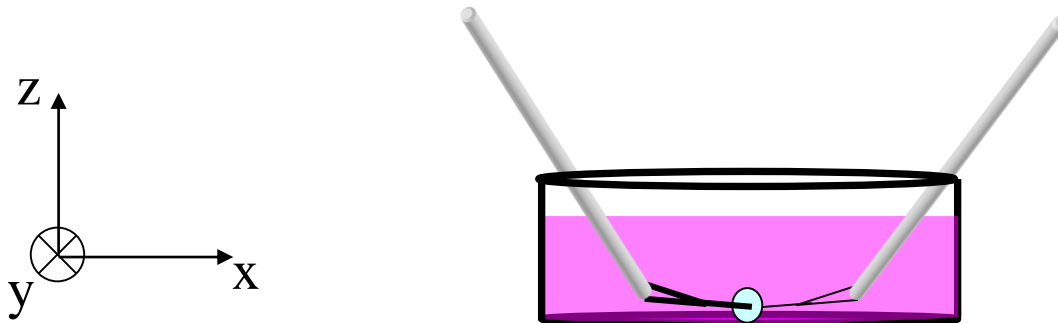
$$\delta = 10 \mu m$$

$$0, 03 nN \cdot \mu m^{-1} < k < 15 nN \cdot \mu m^{-1}$$

$$300 pN < F < 1 \mu N$$

Single cell rheometer

Desprat et al., Rev.Sci. Instrum. 77, 055111-1 (2006)

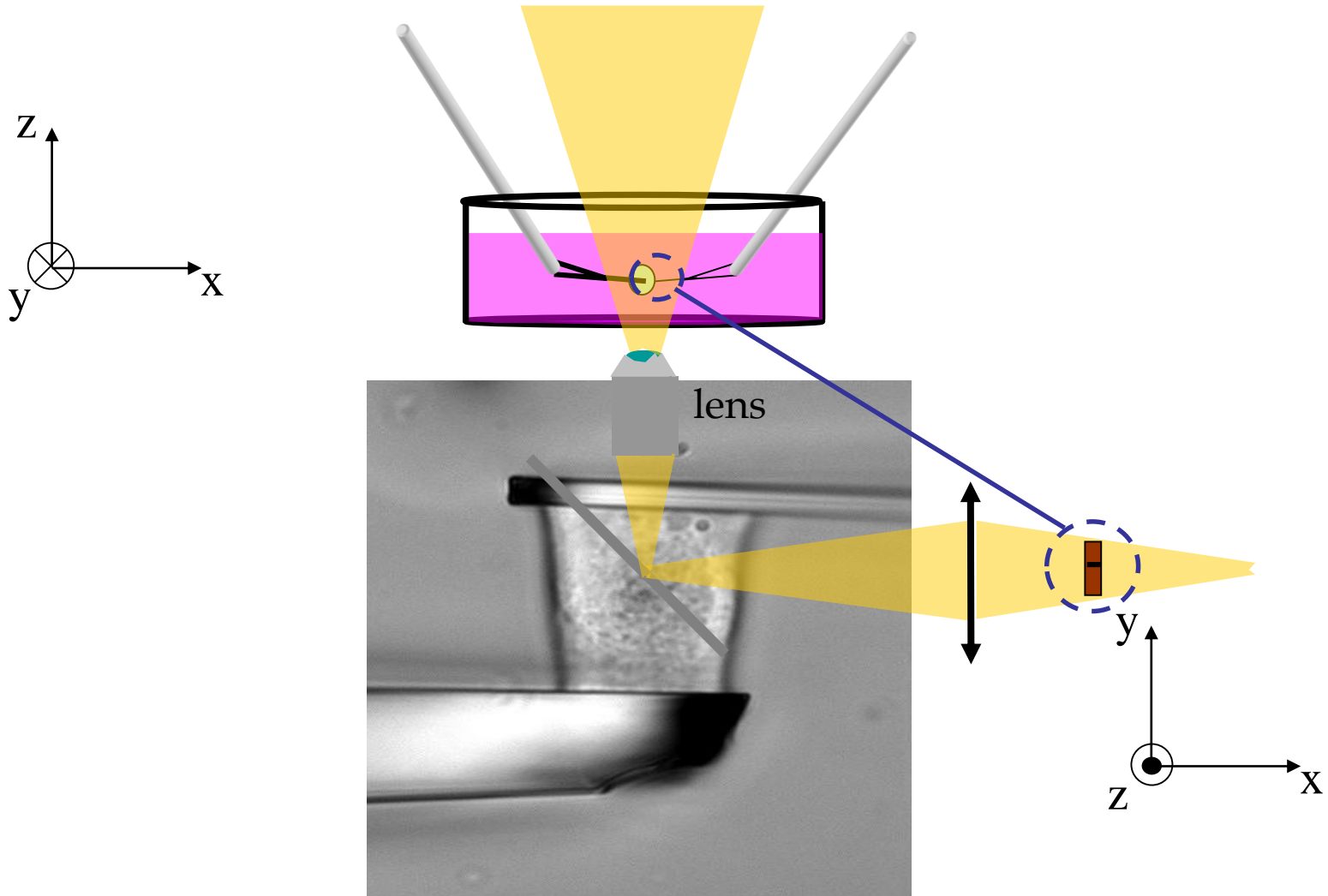


Stretched
plate

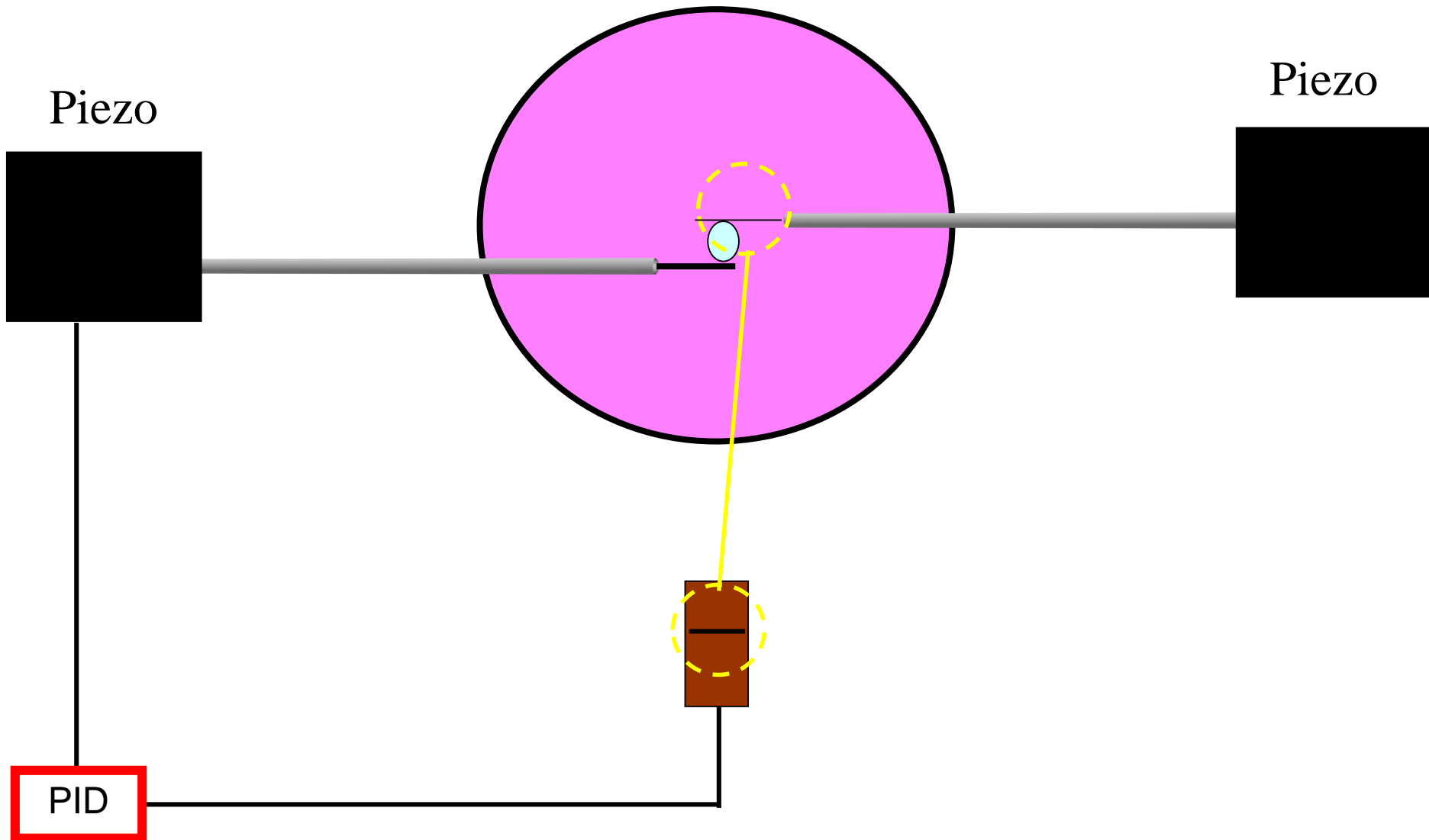
Rigid tube

Single cell rheometer

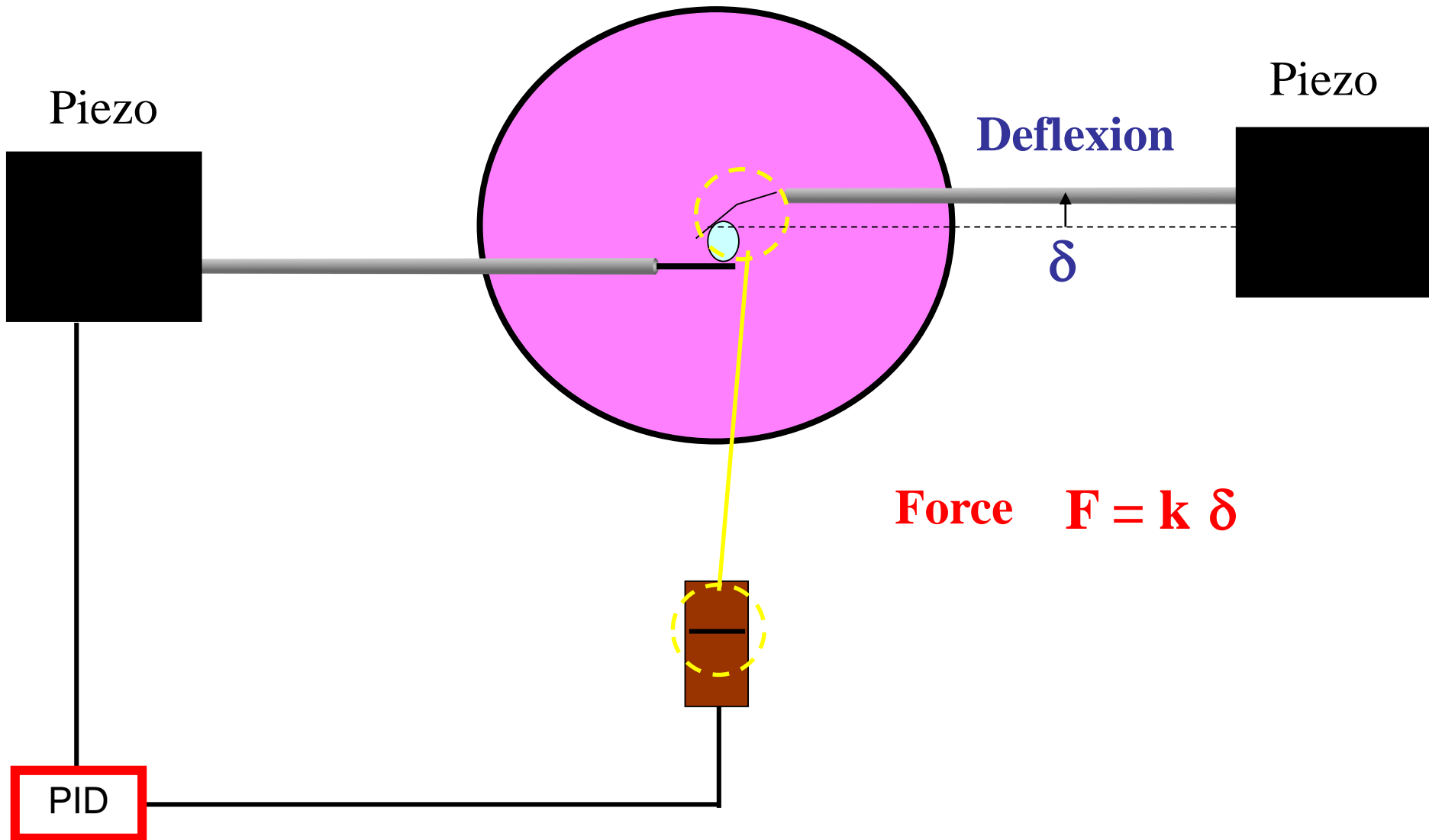
Desprat et al., Rev.Sci. Instrum. 77, 055111-1 (2006)



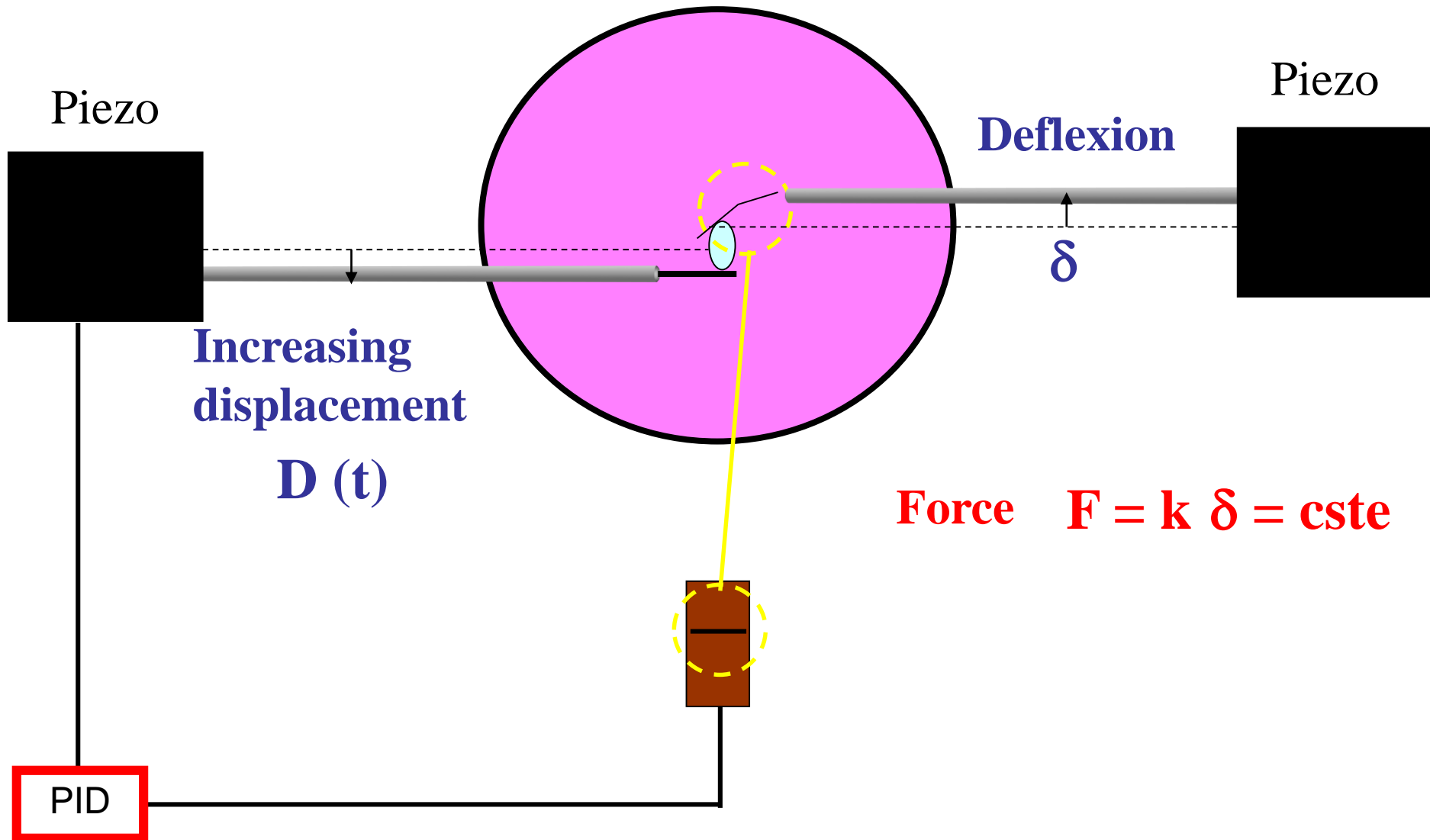
Creep experiment



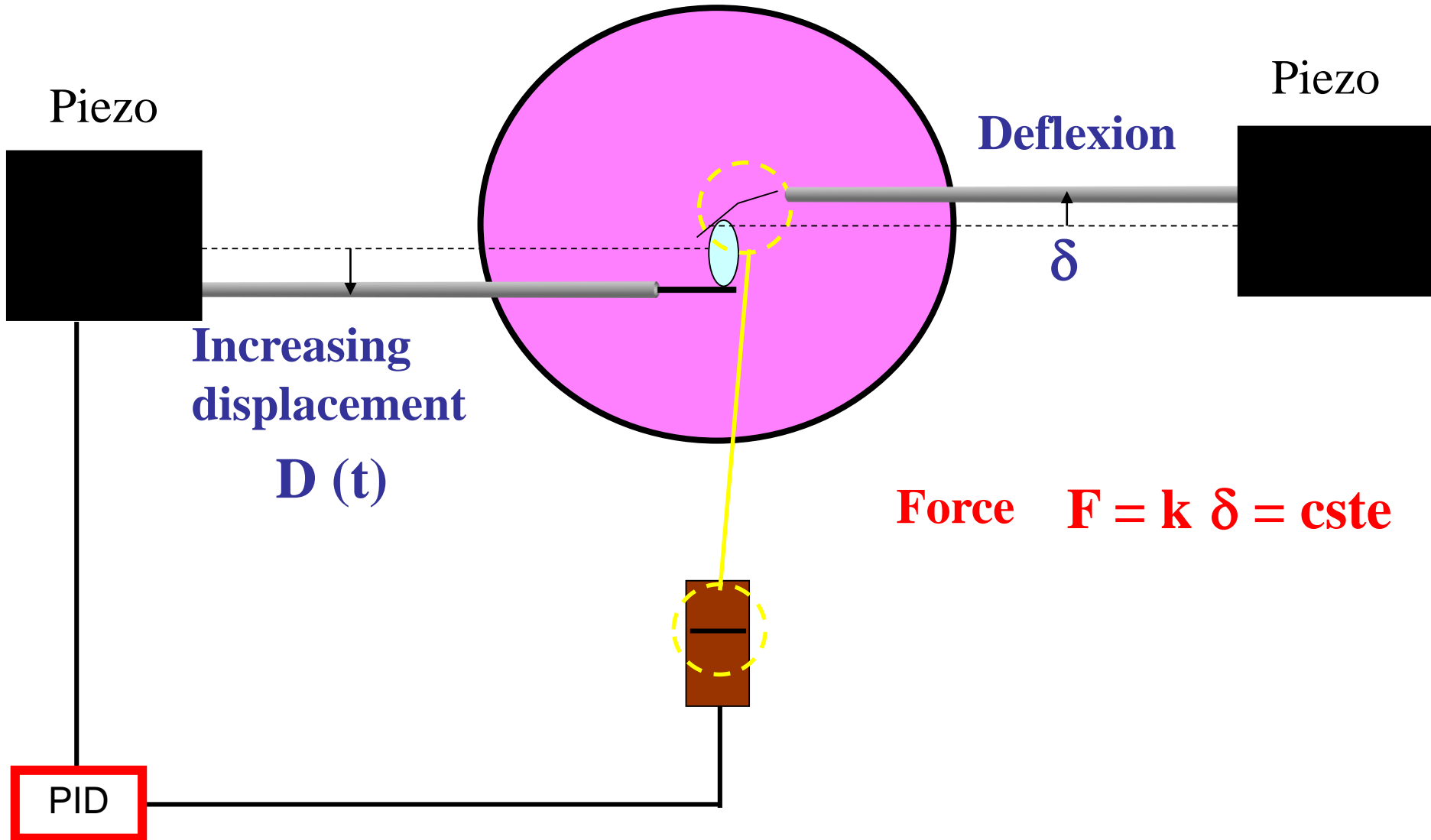
Creep experiment



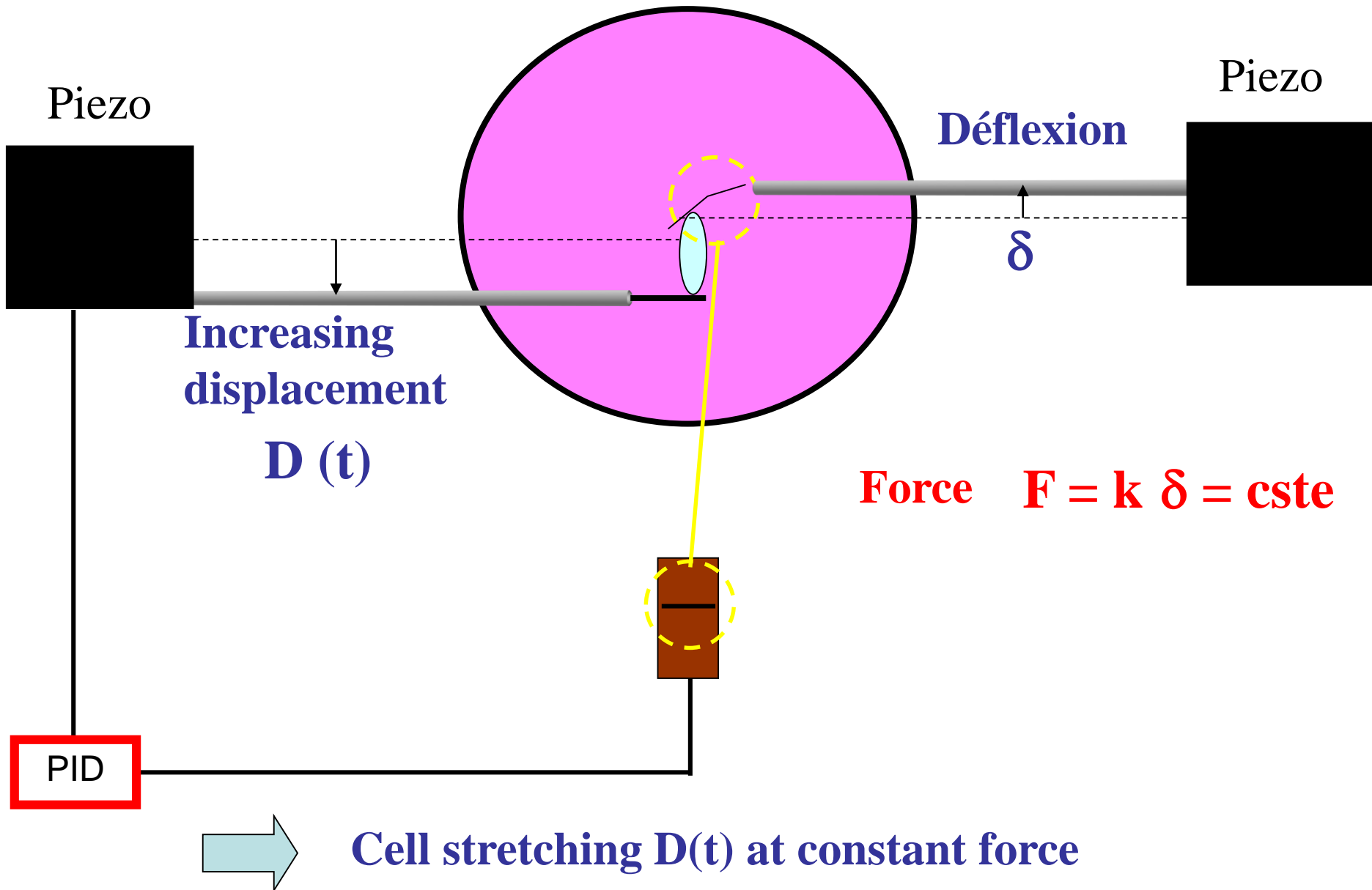
Creep experiment

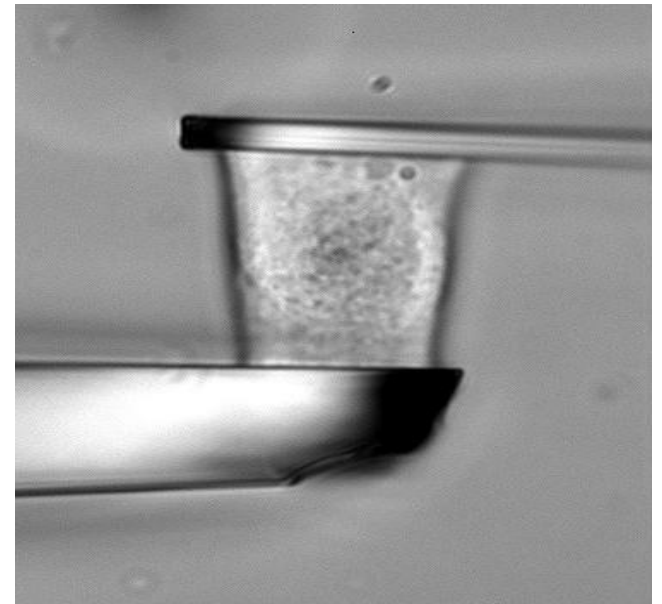
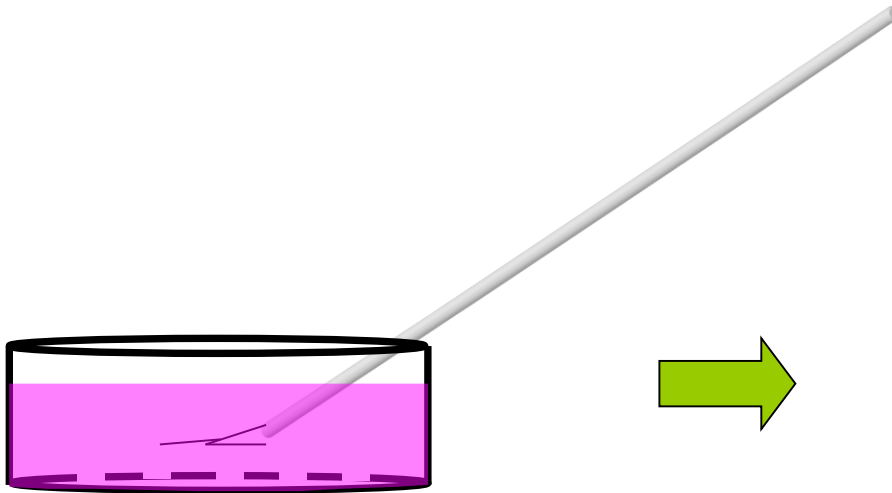
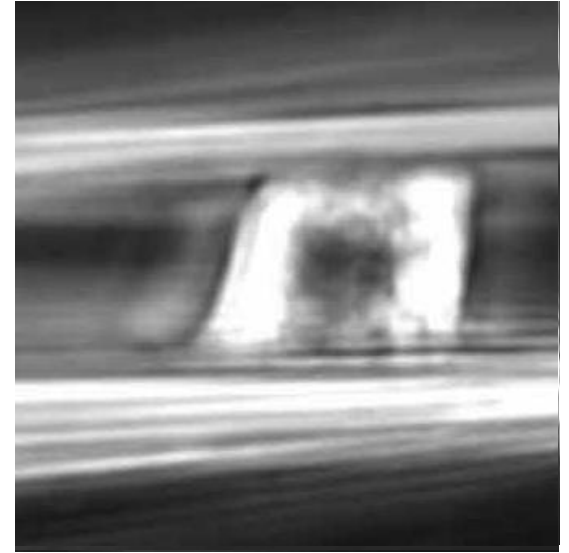
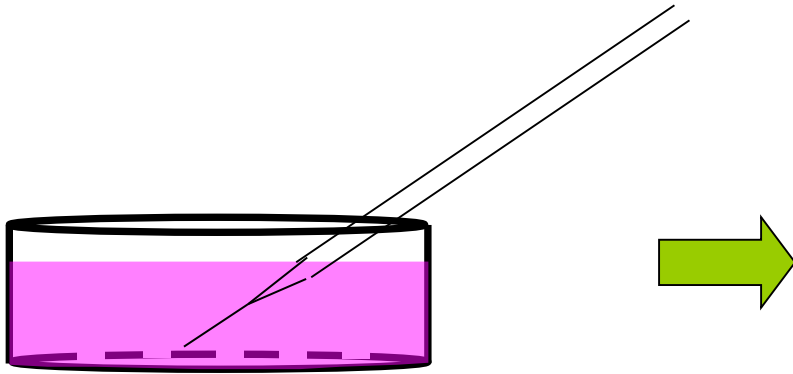


Creep experiment

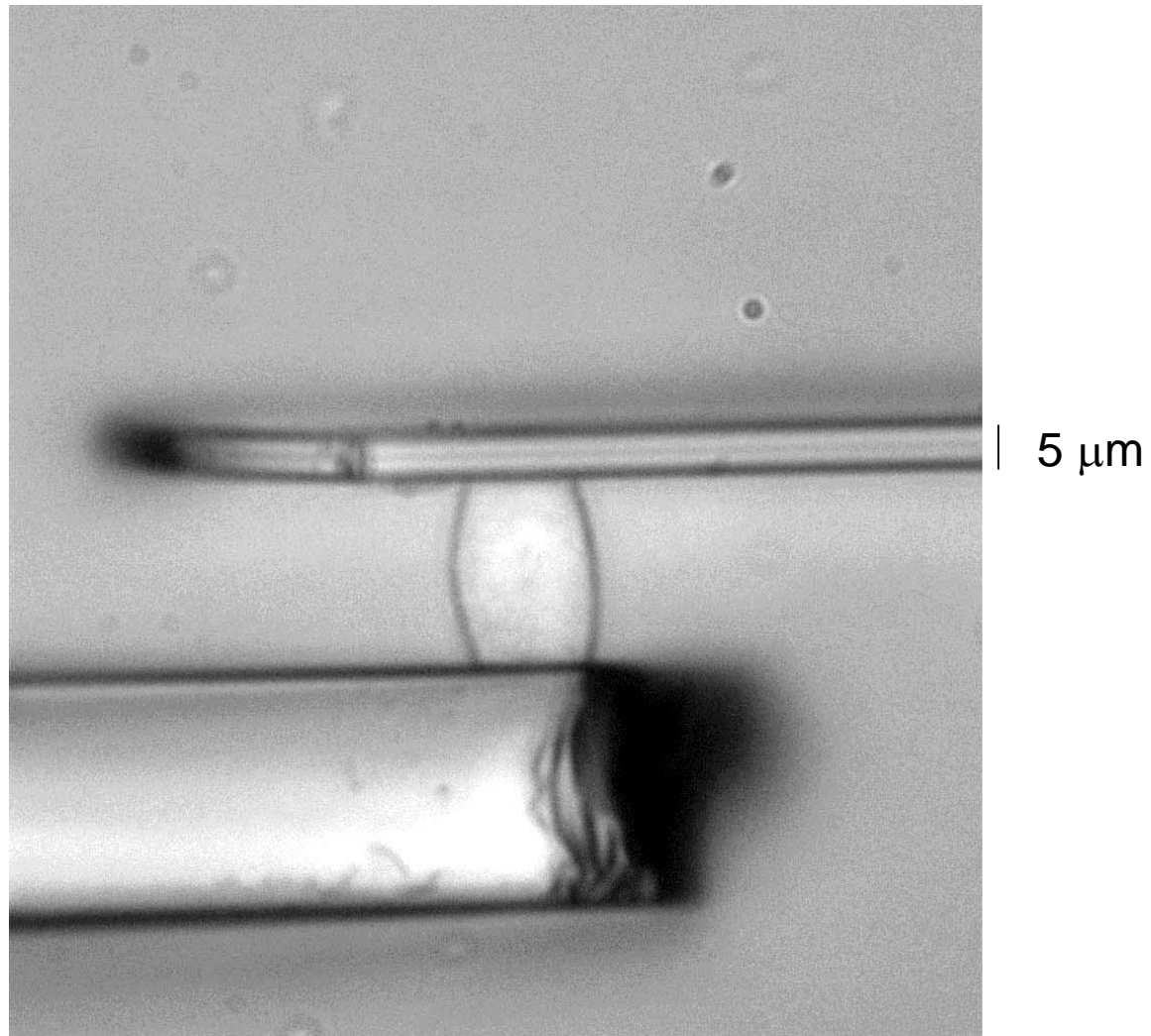


Creep experiment



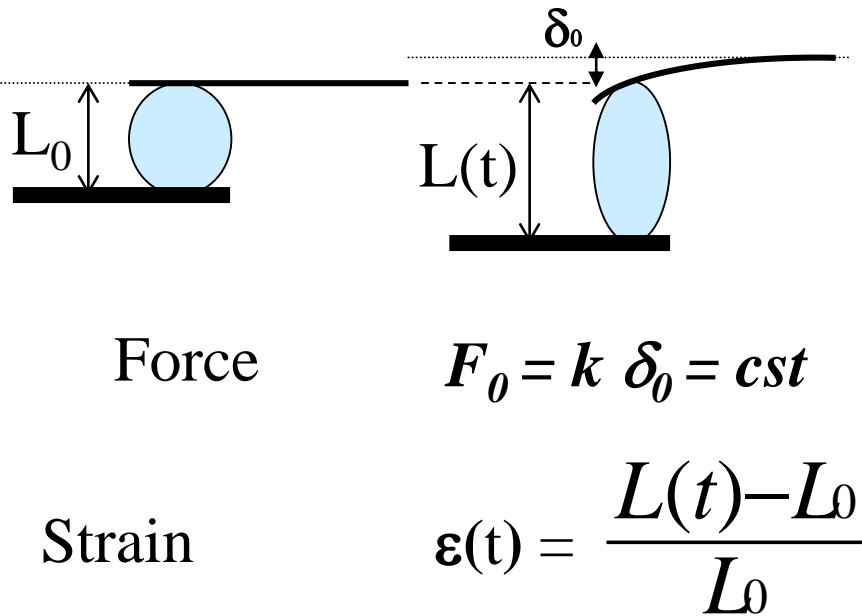
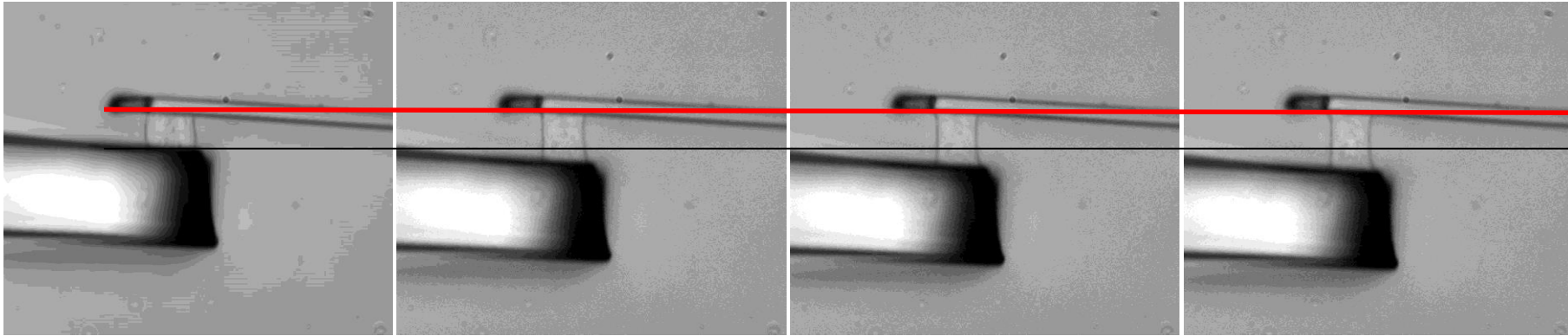


Creep experiment

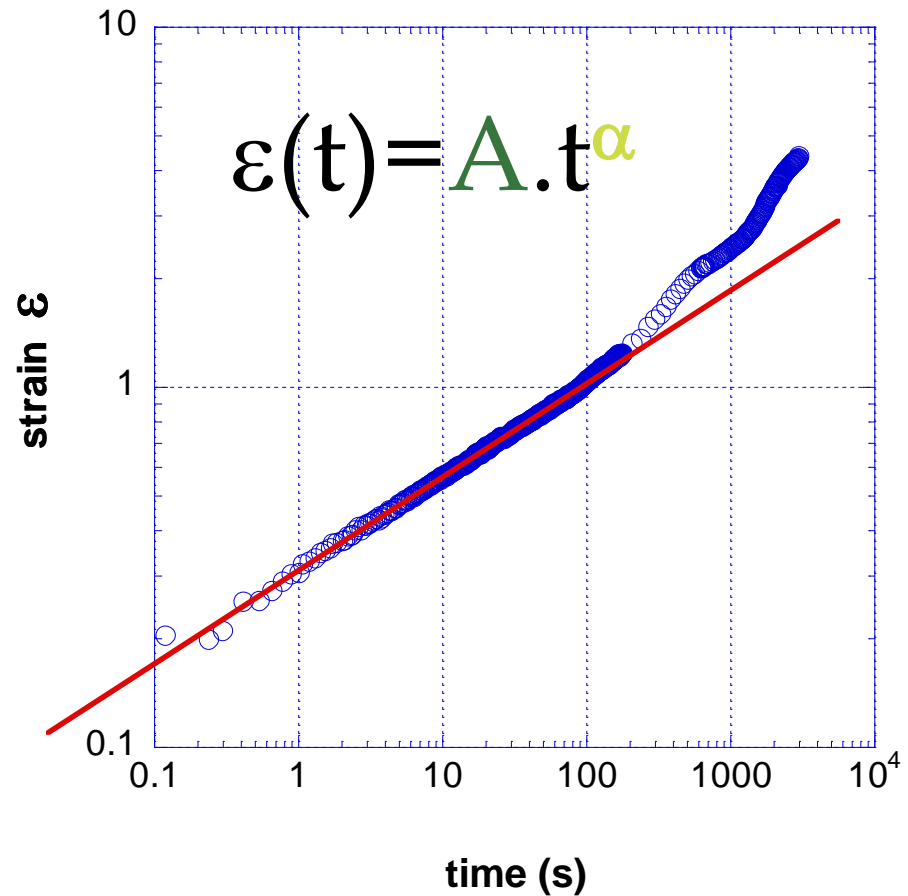


Plates treated with Glutaraldehyde, non specific adhesion

Creep experiment

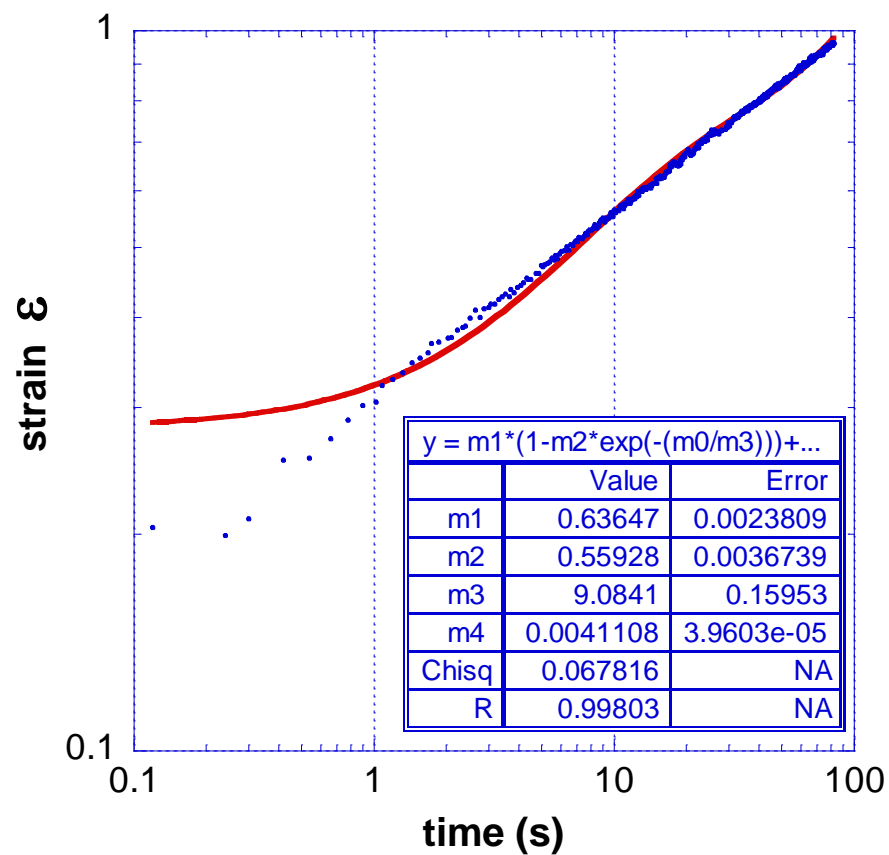
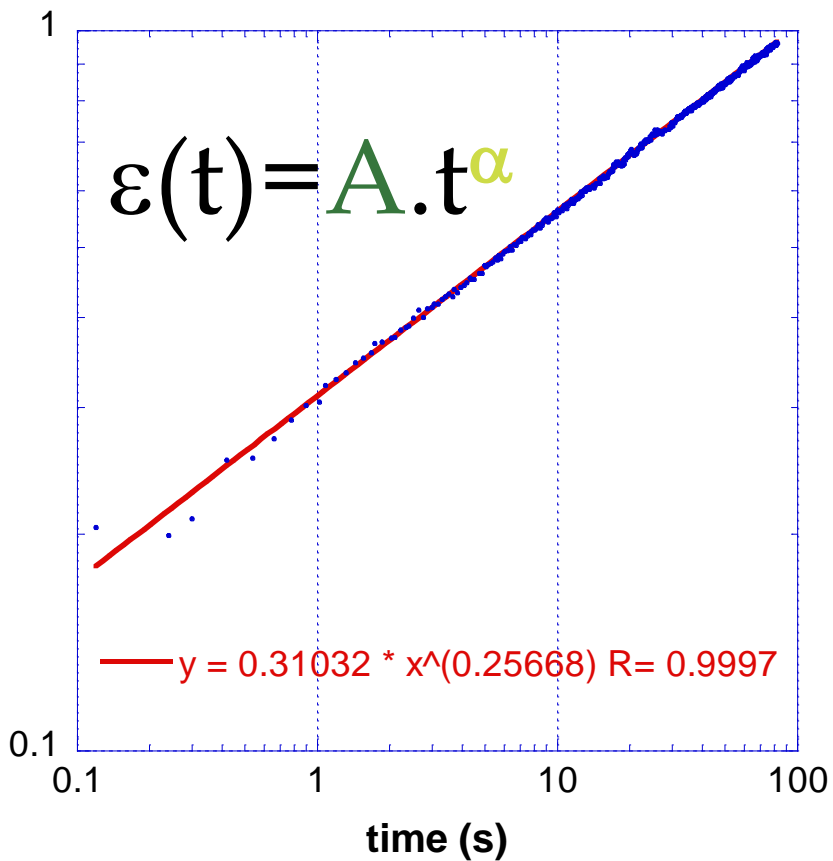
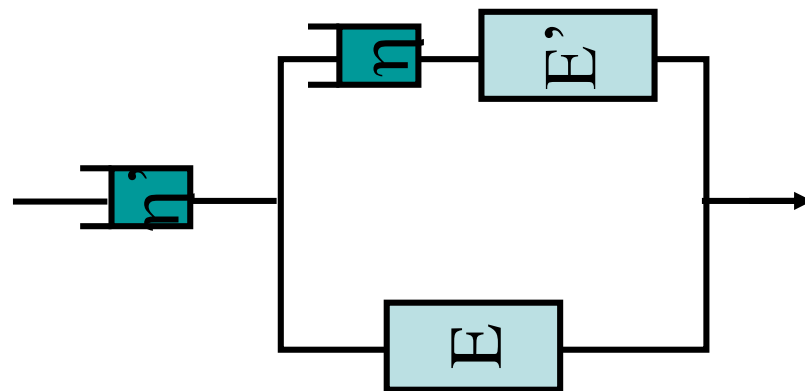


No characteristic time

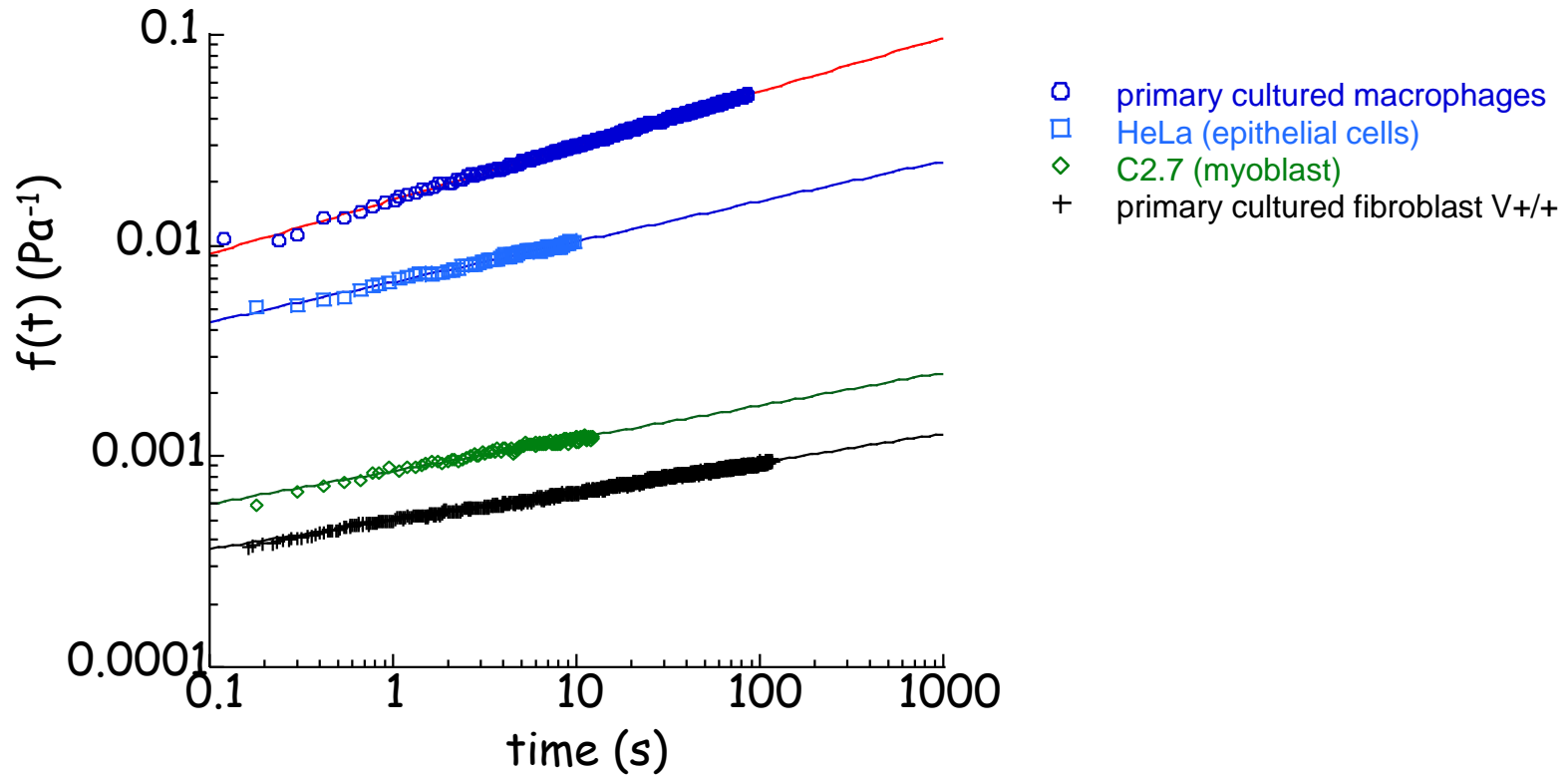


Comparison with simple mechanical models

No characteristic time



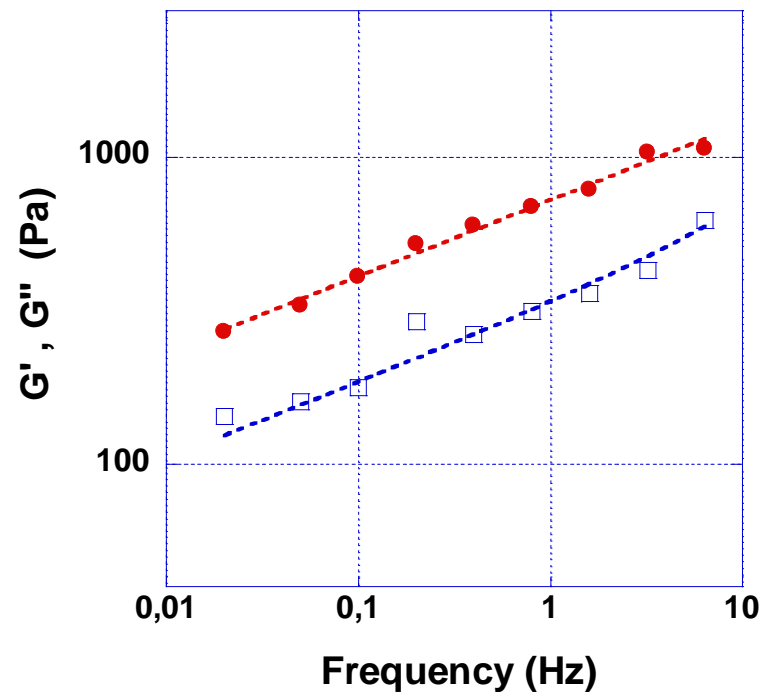
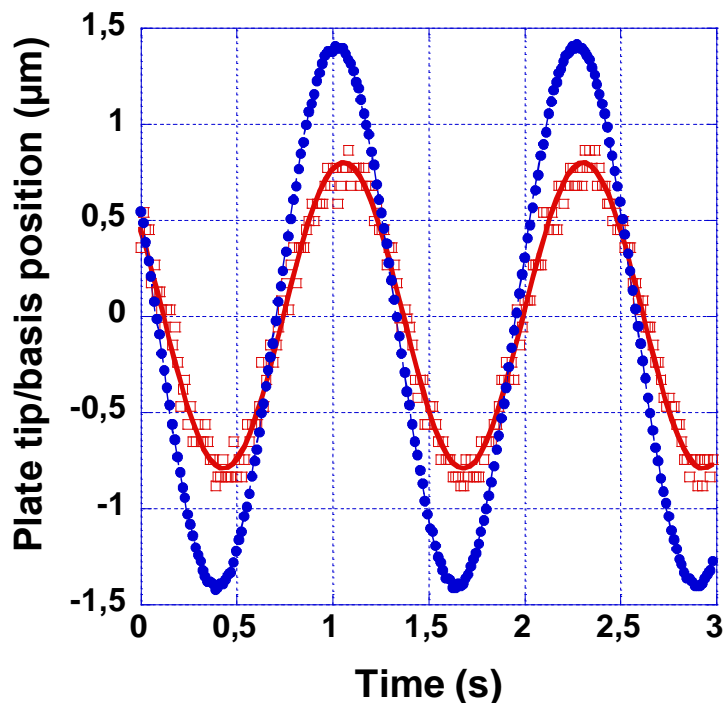
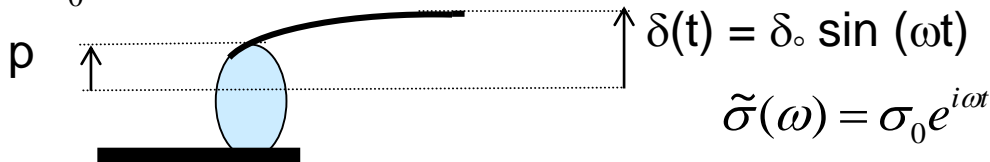
« Universal » behavior



cancer cells F9, J774 alveolar macrophages, A549 alveolar epithelial cells, BEAS-2B of bronchi, human neutrophils

Viscoelastic modulus at small strains

$$\tilde{\varepsilon}(\omega) = \varepsilon_0 e^{i(\omega t + \varphi)}$$



Power law behavior is consistent
Linearity at large strains

Derivations

The fundamental relation
of **linear viscoelasticity**

$$\varepsilon(t) = J(t)\sigma(0) + \int_0^t J(t-t')\dot{\sigma}(t')dt'$$

Then becomes

$$\varepsilon(t) = J(t)\sigma(0) + \int_0^t J(t-t')\sigma(0)\dot{\varepsilon}(t')dt'$$

Laplace transform then yields

$$F(p) = \mathcal{L}\{f(t)\} = \int_{0^-}^{+\infty} e^{-pt} f(t) dt.$$

$$\tilde{\varepsilon}(s) = \frac{\sigma(0)\tilde{J}(s)}{[1 - s\sigma(0)\tilde{J}(s)]}$$

Assuming that $J(t) = At^\alpha$ **as measured in the creep regime**, one finds

$$\varepsilon(t) = \sum_{n=1}^{+\infty} \frac{[\Gamma(1+\alpha)\sigma(0)At^\alpha]^n}{\Gamma(1+n\alpha)}$$

Thus, at high strains, deformation should well be described by a sum of integer powers of the creep function $J(t)$

Soft Glassy Material or ... Fractal Gel

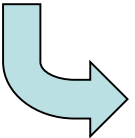
AFM: $L \sim 30$ nm $\alpha \sim 0,20$; $G_0 \sim 710$ Pa (*Alcaraz et al., Biophys J., 2003*)

MTC; OT: $L \sim 3$ μ m $\alpha \sim 0,20$; $G_0 \sim 300$ à 3000 Pa (*Fabry et al., Phys Rev E., 2003*)

(*Balland et al., E. Biophys. J., 2005*)

In agreement with measurements at the cellular scale $L \sim 30$ μ m

$$J(t) = A.t^\alpha \xrightarrow{\text{T.F.}} G'(f) = \frac{(2\pi)^\alpha \cos(\alpha \frac{\pi}{2})}{\Gamma(1+\alpha)} f^{-\alpha}$$

 $G_0 = 660$ Pa

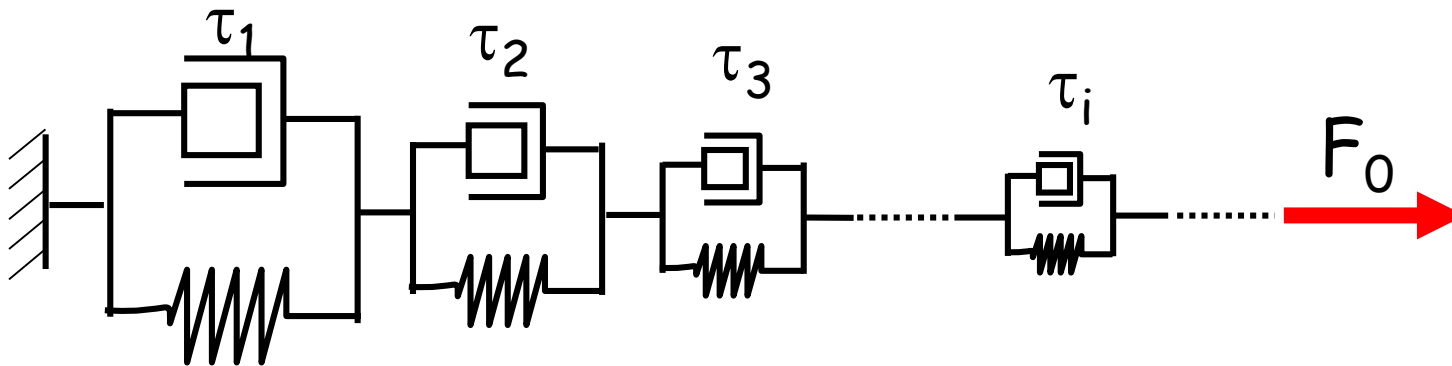
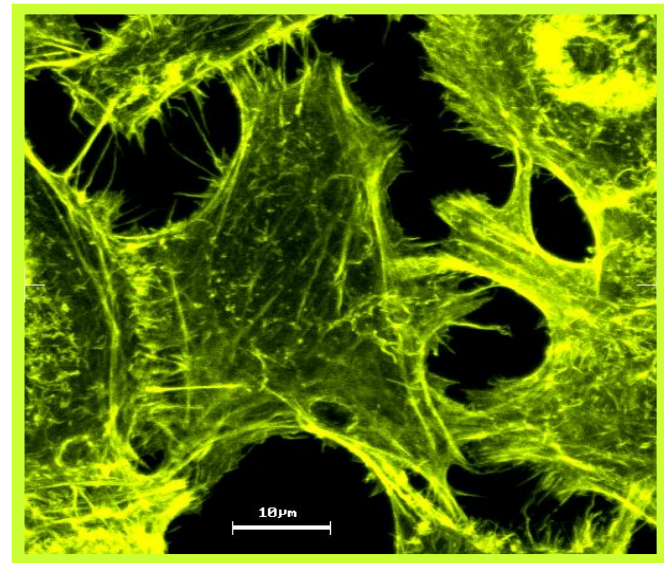
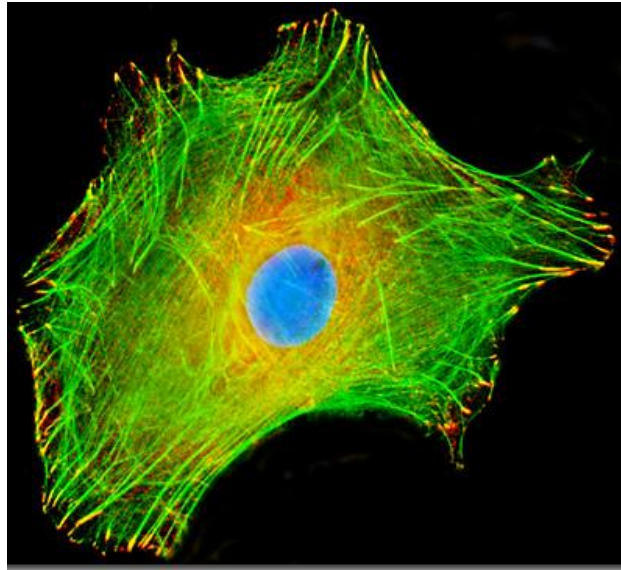
Auto-similarity ?

A simple constitutive model

Actine network:

- individual filaments
- bundles
- fibers

unevenly distributed in
the cell body



The actin network is modeled by an infinite series of nested elementary viscoelastic units with a wide distribution $p(\tau)$ relaxation times τ

Distribution of response times

Balland et al., Phys.Rev.E 74, 021911 (2006)

Simple assumptions:

- $N(d)$ number of units of size d
 $N(d) \sim d^{-a}$ if self similar structure
- relaxation time linked to spatial scale: $\tau \sim d^b$

Then $p(\tau) \sim \tau^{\alpha-2}$ with $\alpha = 1 - a/b$

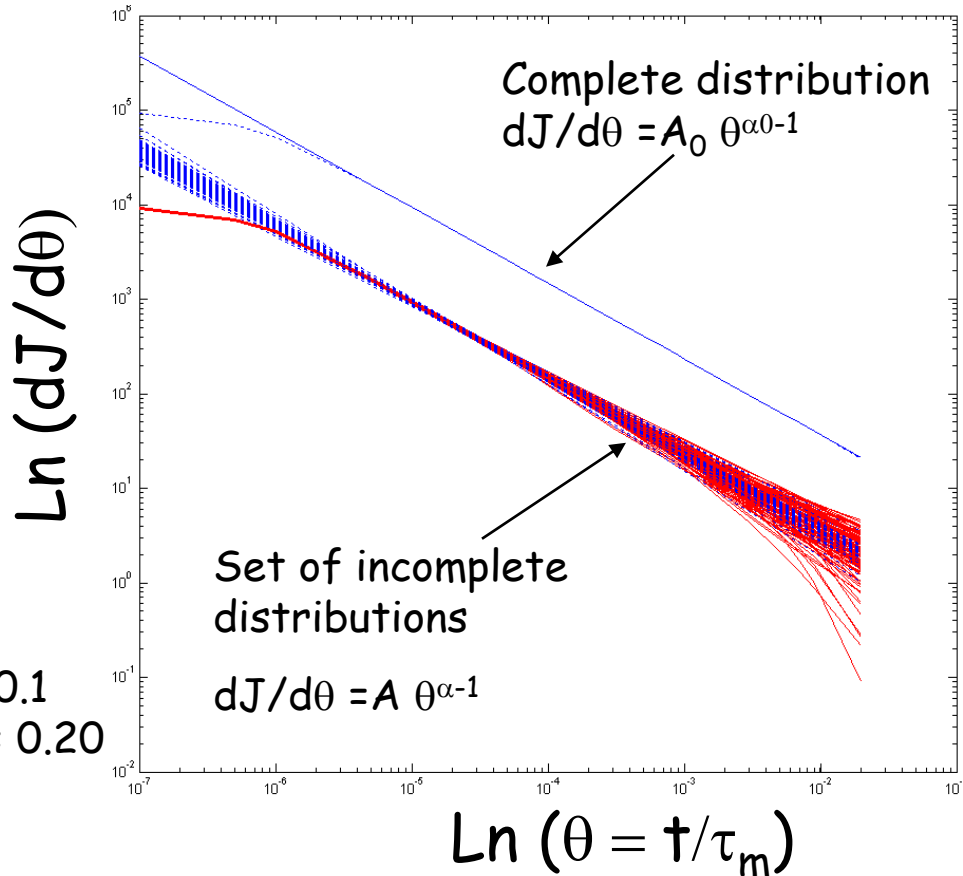
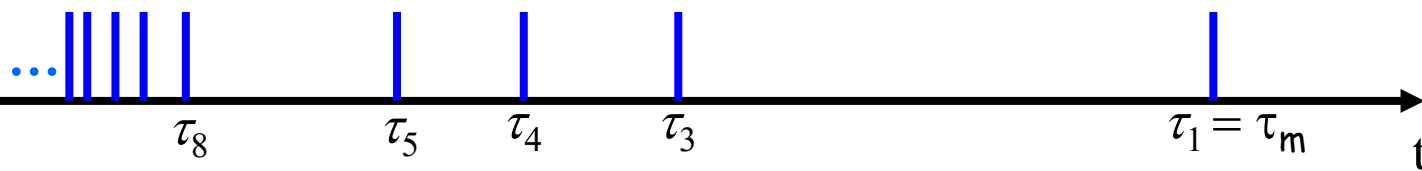
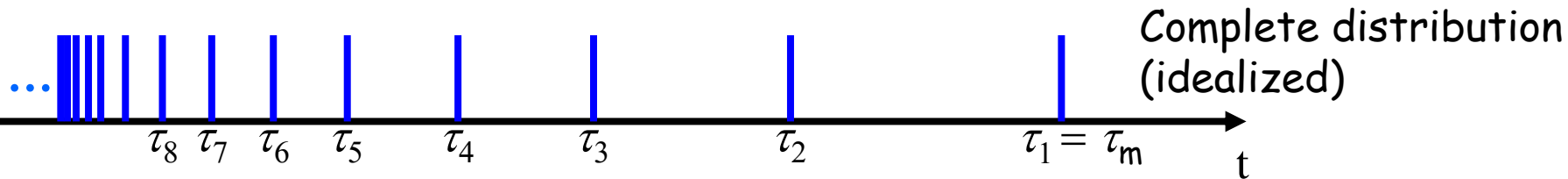
$p(\tau) \sim \tau^{\alpha-2}$ in power law \longrightarrow creep function $J(t)$ as well

$$\frac{dJ}{dt} = \sum_{i=1}^{\infty} \exp\left(-\frac{t}{\tau_i}\right) \approx \int_0^{\infty} \tau^{\alpha-2} \exp\left(-\frac{t}{\tau}\right) d\tau \propto t^{\alpha-1}$$

so $J(t) \sim t^{\alpha}$

Agreement with experimental observations

Response of the system



Incomplete distribution:
 we randomly keep a
 fraction s of the elements

(Simulates the variability
 from one cell to another)

$$dJ/d\theta = A \theta^{\alpha-1}$$

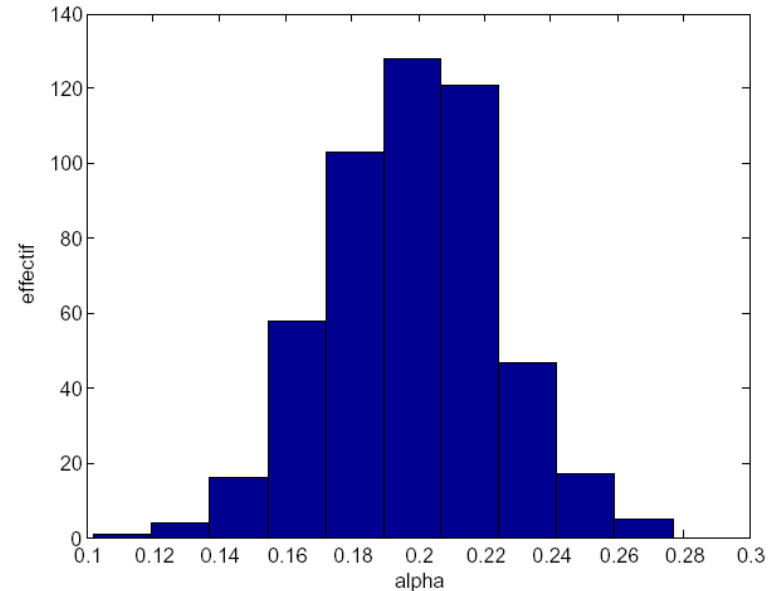
$$J(\theta) = (A/\alpha) \theta^\alpha$$

Dispersion of coefficients of the power law

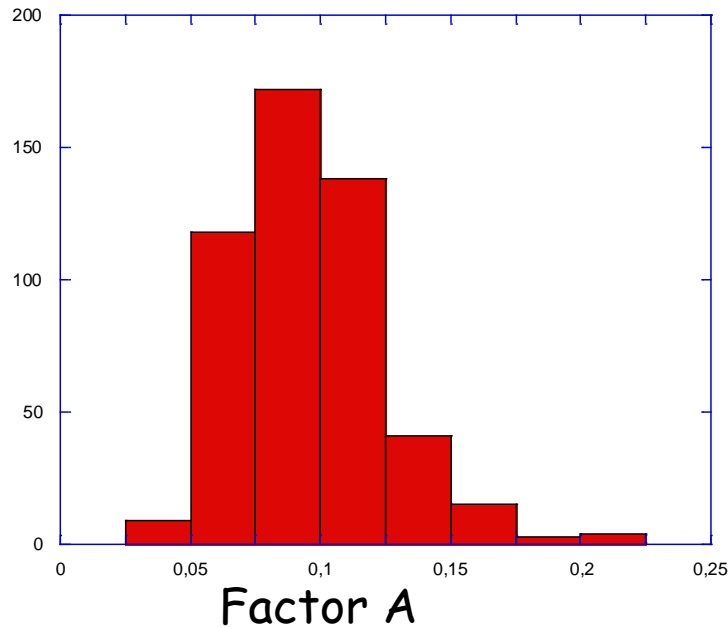
$$J(\theta) = A \theta^\alpha$$

- Normal distribution of exponents α
- Log-normale distribution of factors A

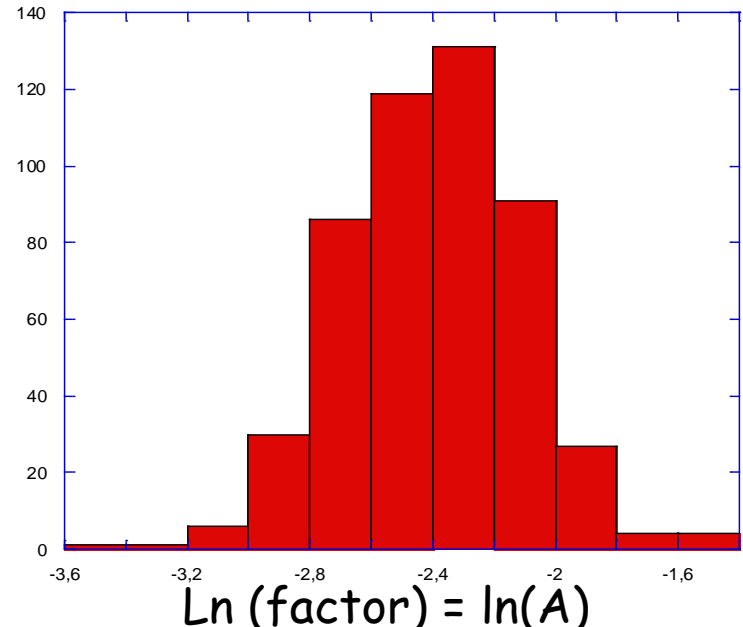
cf experimental results



Exponent α



Factor A



Ln (factor) = $\ln(A)$

Soft Glassy Material or ...

Like foams, emulsion, slurries

Disordered medium with a great number of elements and **out of equilibrium**

Interaction between mesoscopic elements leads to

- large distributions of sizes and relaxation times:
no characteristic time scale
- specific relaxation processes :
non viscous dissipation

Parameter of control **x (noise temperature)**

- **power law** rheological behaviour, **$\alpha = x - 1$**

Possible origins of the power law behavior

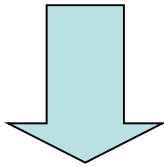
foams, emulsions, pastes, slurries

- Out of equilibrium
- Permanent structural rearrangement



Soft Glassy Materials (SGM)

Sollich, *Phys. Rev. E* (1998)



Dynamic origin

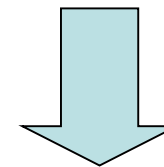
Partially polymerized gels

- Fixed structure
- Fractal dimension



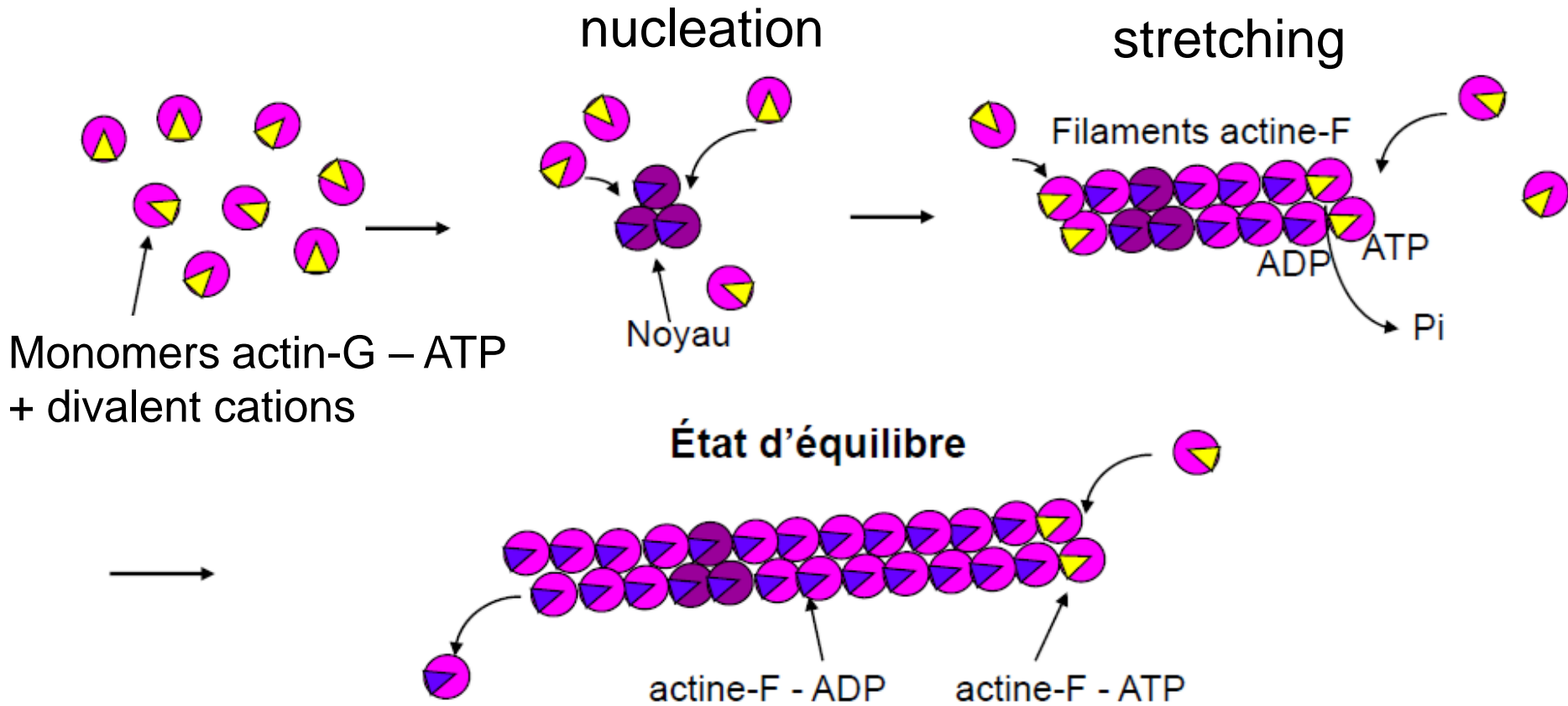
Materials at the « Sol-Gel » transition

Winter *et al.*, *J. of Rheology* (1986)



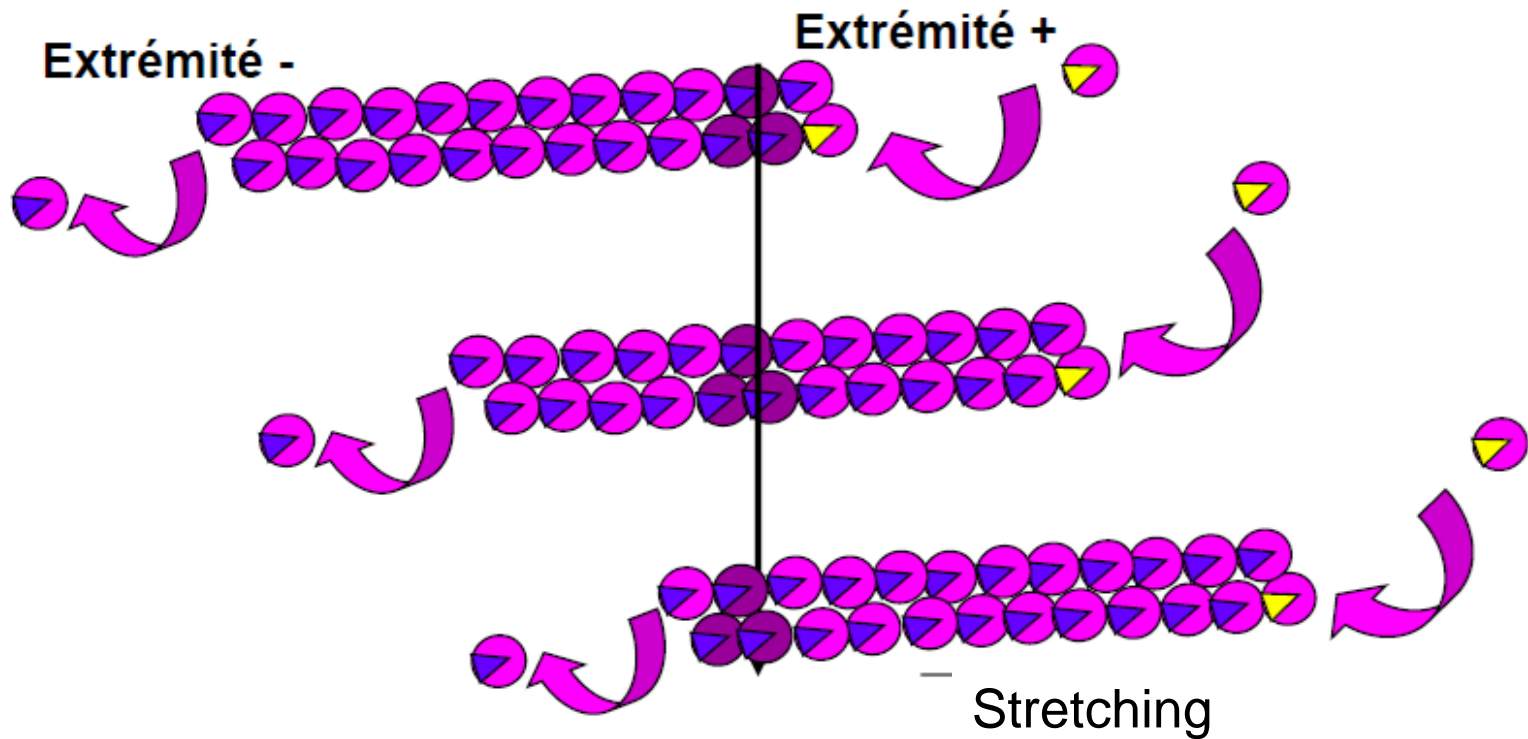
Structural origin

POLYMERIZATION OF ACTIN FILAMENTS

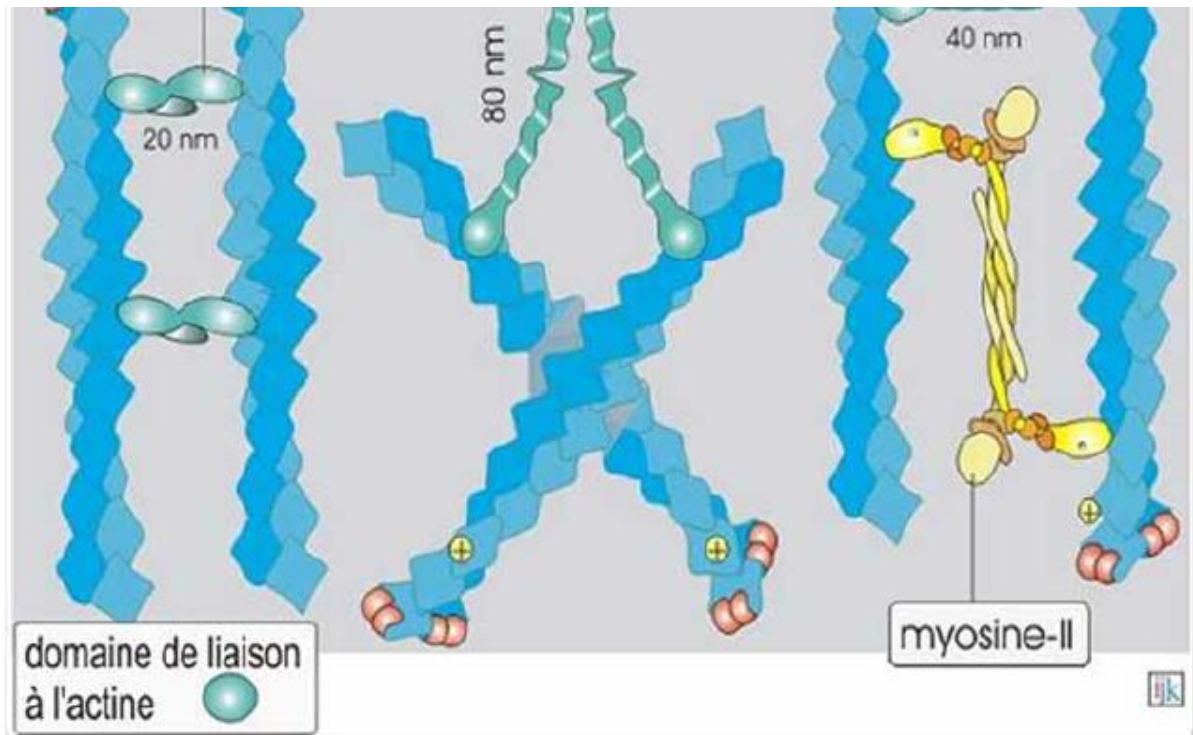


treadmilling

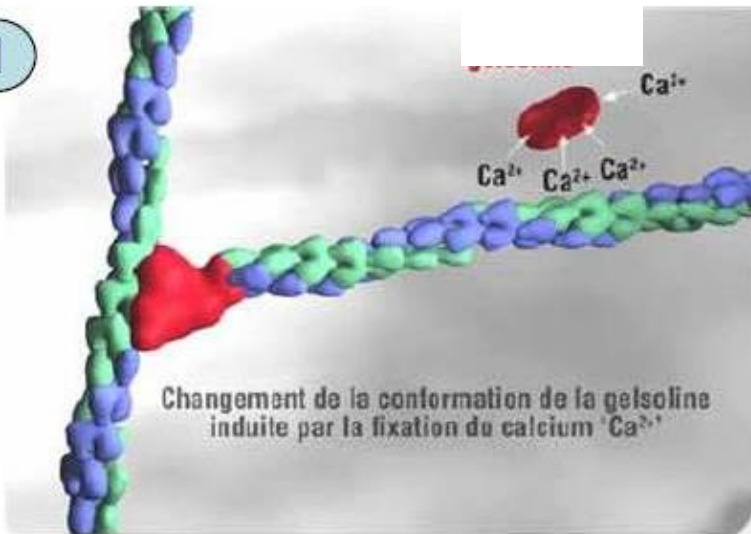
$C_{c-} (0,8\mu M) > C_{actine-G} > C_{c+} (0,1\mu M)$



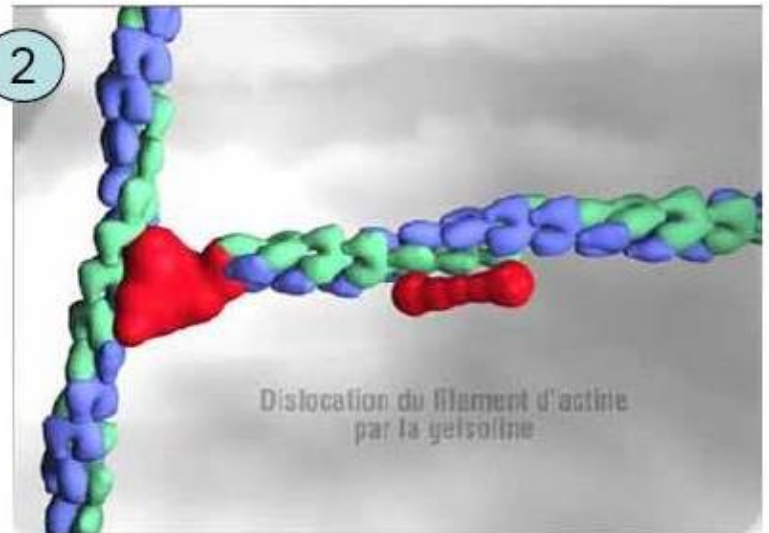
STRUCTURE OF ACTIN FILAMENTS IN THE CELL



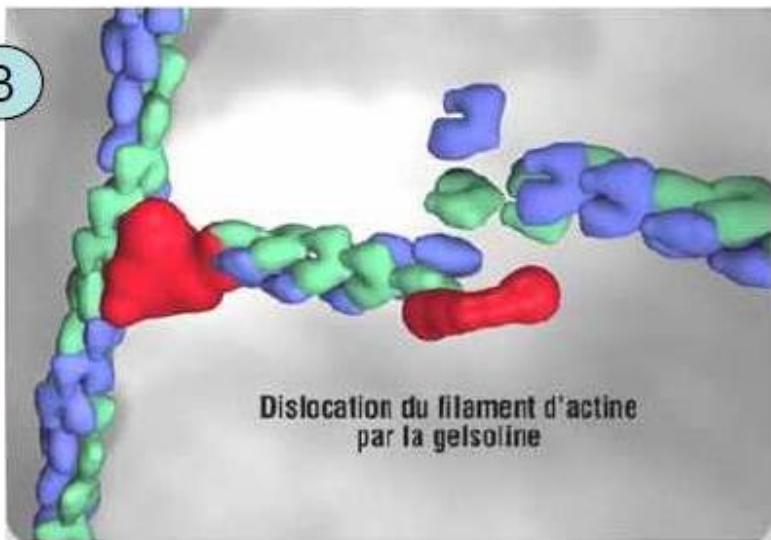
1



2



3



4

